# Physics and geometry of knots-quivers correspondence

Piotr Kucharski

Uppsala University, Sweden

September 27, 2018

< ロ > < 同 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ >

## Collaboration

This talk is based on [1810.xxxx] with:



Pietro Longhi



Tobias Ekholm

with crucial references to [1707.02991, 1707.0417] with:



Markus Reineke



Marko Stošić



Piotr Sułkowski

# Outline



- Knots
- Quivers
- Knots-quivers correspondence
- 2 Physics and geometry of KQ correspondence
  - General idea
  - Physics
  - Geometry

#### Summary

< - 12 ▶

→ Ξ →

- E - N

Knots Quivers Knots-quivers correspondence

# Outline



#### Knots

- Quivers
- Knots-quivers correspondence

#### 2 Physics and geometry of KQ correspondence

- General idea
- Physics
- Geometry

#### 3 Summary

(日) (同) (三) (三)

# Setting the stage

• Ooguri-Vafa M-theory system:



P. Kucharski

Physics and geometry of KQ correspondence

Knots Quivers Knots-quivers correspondence

# Calabi-Yau side (resolved conifold)

- Topological strings ↔ Chern-Simons theory ↔
   ↔ Knot theory [Witten]
- Partition function=HOMFLY-PT gen. series

$$Z = P^{K}(x, a, q) = \sum_{r=0}^{\infty} P_{r}^{K}(a, q) \times^{r}$$

• M5 brane Lagrangian submanifold *L<sub>K</sub>* M2 branes Holomorphic curves



イロト イポト イラト イ

Knots Quivers Knots-quivers correspondence

# Calabi-Yau side (resolved conifold)

- Topological strings ↔ Chern-Simons theory ↔
   ↔ Knot theory [Witten]
- Partition function=HOMFLY-PT gen. series

$$Z = P^{K}(x, a, q) = \sum_{r=0}^{\infty} P_{r}^{K}(a, q) x^{r}$$

• M5 brane Lagrangian submanifold *L<sub>K</sub>* M2 branes Holomorphic curves



(日) (品) (日) (

Knots Quivers Knots-quivers correspondence

# Calabi-Yau side (resolved conifold)

- Topological strings ↔ Chern-Simons theory ↔
   ↔ Knot theory [Witten]
- Partition function=HOMFLY-PT gen. series

$$Z = P^{K}(x, a, q) = \sum_{r=0}^{\infty} P_{r}^{K}(a, q) x^{r}$$

• M5 brane Lagrangian submanifold L<sub>K</sub> M2 branes Holomorphic curves



・ロト ・同ト ・ヨト ・

Flat side  $(\mathbb{R}^{1,4})$ 

- $T[L_K]$ : 3d  $\mathcal{N} = 2$  eff. theory on  $\mathbb{R}^{1,2}$
- Partition function=generating function of Labastida-Mariño-Ooguri-Vafa invariants N<sup>K</sup><sub>r,i,j</sub>

$$Z = \operatorname{Exp}\left(\frac{\sum_{r,i,j} N_{r,i,j}^{K} x^{r} a^{i} q^{j}}{1 - q^{2}}\right)$$

•  $\frac{M5 \text{ brane}}{M2 \text{ branes}} = \frac{\mathbb{R}^{1,2}}{\mathbb{BPS} \text{ states}}$ 

 $\bullet~$  BPS states are counted by LMOV invariants  $\Rightarrow~$  LMOV integrality conjecture



(日) (品) (日) (

Flat side  $(\mathbb{R}^{1,4})$ 

- $T[L_K]$ : 3d  $\mathcal{N} = 2$  eff. theory on  $\mathbb{R}^{1,2}$
- Partition function=generating function of Labastida-Mariño-Ooguri-Vafa invariants N<sup>K</sup><sub>r,i,i</sub>

$$Z = \operatorname{Exp}\left(\frac{\sum_{r,i,j} N_{r,i,j}^{K} x^{r} a^{i} q^{j}}{1 - q^{2}}\right)$$

•  $\frac{M5 \text{ brane}}{M2 \text{ branes}} = \frac{\mathbb{R}^{1,2}}{BPS \text{ states}}$ 

 $\bullet~$  BPS states are counted by LMOV invariants  $\Rightarrow~$  LMOV integrality conjecture



イロト イポト イラト イラ

Flat side  $(\mathbb{R}^{1,4})$ 

- $T[L_K]$ : 3d  $\mathcal{N} = 2$  eff. theory on  $\mathbb{R}^{1,2}$
- Partition function=generating function of Labastida-Mariño-Ooguri-Vafa invariants N<sup>K</sup><sub>r,i,i</sub>

$$Z = \operatorname{Exp}\left(\frac{\sum_{r,i,j} N_{r,i,j}^{K} x^{r} a^{i} q^{j}}{1 - q^{2}}\right)$$

• 
$$\frac{M5 \text{ brane}}{M2 \text{ branes}} = \frac{\mathbb{R}^{1,2}}{\mathbb{R}^{1,2}}$$

 $\bullet~$  BPS states are counted by LMOV invariants  $\Rightarrow~$  LMOV integrality conjecture



イロト イポト イヨト イヨ

Flat side  $(\mathbb{R}^{1,4})$ 

- $T[L_K]$ : 3d  $\mathcal{N} = 2$  eff. theory on  $\mathbb{R}^{1,2}$
- Partition function=generating function of Labastida-Mariño-Ooguri-Vafa invariants N<sup>K</sup><sub>r,i,i</sub>

$$Z = \operatorname{Exp}\left(\frac{\sum_{r,i,j} N_{r,i,j}^{K} x^{r} a^{i} q^{j}}{1 - q^{2}}\right)$$

•  $\frac{M5 \text{ brane}}{M2 \text{ branes}} \frac{\mathbb{R}^{1,2}}{BPS \text{ states}}$ 

 $\bullet~$  BPS states are counted by LMOV invariants  $\Rightarrow~$  LMOV integrality conjecture

・ 同 ト ・ ラ ト ・



Knots Quivers Knots-quivers correspondence

# Outline



- Knots
- Quivers
- Knots-quivers correspondence
- 2 Physics and geometry of KQ correspondence
  - General idea
  - Physics
  - Geometry

#### 3 Summary

(日) (同) (三) (三)

## Quivers

- Quiver  $Q = (Q_0, Q_1)$ 
  - Q<sub>0</sub> is a set of vertices
  - $Q_1$  is a set of arrows between them (loops allowed)
- Matrix form: *C<sub>ij</sub>* gives the number of arrows between vertices *i* and *j* 
  - If  $C_{ij} = C_{ji}$  we call the quiver symmetric
- Quiver representations:  $(Q_0, Q_1) \longrightarrow ($ Vector spaces, Linear maps
- Example from the picture:
  - $Q_0 = \{1,2\} \longrightarrow$  vector spaces  $V_1$  and  $V_2$  of dimensions  $d_1$  and  $d_2$  (dimension vector  $\boldsymbol{d} = (d_1, d_2)$ )
  - $Q_1 = \{1 \rightarrow 1\} \longrightarrow$  linear map  $V_1 \rightarrow V_1$  , as the set of  $P_1 \rightarrow P_1$

2

 $\left|\begin{array}{cc}1&0\\0&0\end{array}\right|$ 

# Quivers

- Quiver  $Q = (Q_0, Q_1)$ 
  - $Q_0$  is a set of vertices
  - $Q_1$  is a set of arrows between them (loops allowed)
- Matrix form: *C<sub>ij</sub>* gives the number of arrows between vertices *i* and *j* 
  - If  $C_{ij} = C_{ji}$  we call the quiver symmetric
- Quiver representations:  $(Q_0, Q_1) \longrightarrow ($ Vector spaces, Linear maps
- Example from the picture:
  - $Q_0 = \{1,2\} \longrightarrow$  vector spaces  $V_1$  and  $V_2$  of dimensions  $d_1$  and  $d_2$  (dimension vector  $\boldsymbol{d} = (d_1, d_2)$ )
  - $Q_1 = \{1 \to 1\} \longrightarrow$  linear map  $V_1 \to V_1$  , as above the set of  $Q_1 = \{1 \to 1\}$

2

 $\begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$ 

naa

# Quivers

- Quiver  $Q = (Q_0, Q_1)$ 
  - Q<sub>0</sub> is a set of vertices
  - $Q_1$  is a set of arrows between them (loops allowed)
- Matrix form: *C<sub>ij</sub>* gives the number of arrows between vertices *i* and *j* 
  - If  $C_{ij} = C_{ji}$  we call the quiver symmetric
- Quiver representations:  $(Q_0, Q_1) \longrightarrow ($ Vector spaces, Linear maps)
- Example from the picture:
  - $Q_0 = \{1,2\} \longrightarrow$  vector spaces  $V_1$  and  $V_2$  of dimensions  $d_1$  and  $d_2$  (dimension vector  $\boldsymbol{d} = (d_1, d_2)$ )
  - $Q_1 = \{1 \to 1\} \longrightarrow$  linear map  $V_1 \to V_1$  , as above the set of  $Q_1 = \{1 \to 1\}$

2

 $\left|\begin{array}{cc}1&0\\0&0\end{array}\right|$ 

SQA

# Quivers

- Quiver  $Q = (Q_0, Q_1)$ 
  - Q<sub>0</sub> is a set of vertices
  - $Q_1$  is a set of arrows between them (loops allowed)
- Matrix form: *C<sub>ij</sub>* gives the number of arrows between vertices *i* and *j* 
  - If  $C_{ij} = C_{ji}$  we call the quiver symmetric
- Quiver representations:  $(Q_0, Q_1) \longrightarrow ($ Vector spaces, Linear maps)
- Example from the picture:
  - $Q_0 = \{1,2\} \longrightarrow$  vector spaces  $V_1$  and  $V_2$  of dimensions  $d_1$  and  $d_2$  (dimension vector  $\boldsymbol{d} = (d_1, d_2)$ )
  - $Q_1 = \{1 \rightarrow 1\} \longrightarrow$  linear map  $V_1 \rightarrow V_1$  , as the set of  $P_1 \rightarrow P_1$

2

 $\left|\begin{array}{cc}1&0\\0&0\end{array}\right|$ 

SQA

# Quivers

- Quiver  $Q = (Q_0, Q_1)$ 
  - Q<sub>0</sub> is a set of vertices
  - $Q_1$  is a set of arrows between them (loops allowed)
- Matrix form: *C<sub>ij</sub>* gives the number of arrows between vertices *i* and *j* 
  - If  $C_{ij} = C_{ji}$  we call the quiver symmetric
- Quiver representations:  $(Q_0, Q_1) \longrightarrow ($ Vector spaces, Linear maps)
- Example from the picture:
  - $Q_0 = \{1,2\} \longrightarrow$  vector spaces  $V_1$  and  $V_2$  of dimensions  $d_1$  and  $d_2$  (dimension vector  $\boldsymbol{d} = (d_1, d_2)$ )
  - $Q_1 = \{1 \rightarrow 1\} \longrightarrow \text{linear map } V_1 \rightarrow V_1$

2

 $\left|\begin{array}{cc}1&0\\0&0\end{array}\right|$ 

200

Knots Quivers Knots-quivers correspondence

## Motivic generating series and DT invariants

• Motivic generating series encodes information about moduli spaces of representations of  ${\cal Q}$ 

$$P^{Q}(\mathbf{x},q) = \sum_{d_{1},...,d_{|Q_{0}|} \ge 0} (-q)^{\sum_{1 \le i,j \le |Q_{0}|} C_{i,j}d_{i}d_{j}} \prod_{i=1}^{|Q_{0}|} \frac{x_{i}^{d_{i}}}{(q^{2};q^{2})_{d_{i}}} = Exp\left(\frac{\Omega^{Q}(\mathbf{x},q)}{1-q^{2}}\right)$$

- Ω<sup>Q</sup>(x,q) is a generating function of motivic Donaldson-Thomas (DT) invariants
- Exp is the plethystic exponential

 $\begin{aligned} & \operatorname{Exp}\left(x^{r}a^{j}q^{j}\right) = (1 - x^{r}a^{j}q^{j})^{-1} \\ & \operatorname{Exp}\left(f + g\right) = \operatorname{Exp}\left(f\right)\operatorname{Exp}\left(g\right)_{q \text{ is } q \text{ if } q \text{ is } q \text{ if } q \text{ is } q \text{ if }$ 

 $\begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$ 

Knots Quivers Knots-quivers correspondence

#### Motivic generating series and DT invariants

• Motivic generating series encodes information about moduli spaces of representations of *Q* 

$$P^{Q}(\mathbf{x},q) = \sum_{d_{1},...,d_{|Q_{0}|} \ge 0} (-q)^{\sum_{1 \le i,j \le |Q_{0}|} C_{i,j}d_{i}d_{j}} \prod_{i=1}^{|Q_{0}|} \frac{x_{i}^{d_{i}}}{(q^{2};q^{2})_{d_{i}}} = \exp\left(\frac{\Omega^{Q}(\mathbf{x},q)}{1-q^{2}}\right)$$

- Ω<sup>Q</sup>(x,q) is a generating function of motivic Donaldson-Thomas (DT) invariants
- $(q^2; q^2)_{d_i}$  is the q-Pochhammer symbol

$$(z;q^{2})_{r} := \prod_{i=0}^{r-1} (1 - zq^{2i}) = (1 - z) (1 - zq^{2}) \dots (1 - zq^{2(r-1)})$$

P. Kucharski Physics and geometry of KQ correspondence

 $\begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$ 

Knots Quivers Knots-quivers correspondence

#### Motivic generating series and DT invariants

• Motivic generating series encodes information about moduli spaces of representations of Q

$$P^{Q}(\mathbf{x},q) = \sum_{d_{1},...,d_{|Q_{0}|} \ge 0} (-q)^{\sum_{1 \le i,j \le |Q_{0}|} C_{i,j}d_{i}d_{j}} \prod_{i=1}^{|Q_{0}|} \frac{x_{i}^{d_{i}}}{(q^{2};q^{2})_{d_{i}}} = \exp\left(\frac{\Omega^{Q}(\mathbf{x},q)}{1-q^{2}}\right)$$

• Example from the picture

$$P^{Q}(x_{1}, x_{2}, q) = \sum_{d_{1}, d_{2} \ge 0} (-q)^{d_{1}^{2}} \frac{x_{1}^{d_{1}}}{(q^{2}; q^{2})_{d_{1}}} \frac{x_{2}^{d_{2}}}{(q^{2}; q^{2})_{d_{2}}} \qquad \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$$
$$= \operatorname{Exp}\left(\frac{-qx_{1} + x_{2}}{1 - q^{2}}\right)$$

P. Kucharski Physics and geometry of KQ correspondence

# Outline



- Knots
- Quivers
- Knots-quivers correspondence
- 2 Physics and geometry of KQ correspondence
  - General idea
  - Physics
  - Geometry

#### 3 Summary

イロト イポト イヨト イヨト

Knots Quivers Knots-quivers correspondence

### Knots-quivers correspondence



• Knots-quivers (KQ) correspondence is an equality

$$P^{K}(x,a,q) = P^{Q_{K}}(\mathbf{x},q)\Big|_{x_{i}=a^{a_{i}}q^{q_{i}-C_{ii}}x}$$

•  $P^{K}(x, a, q)$  – HOMFLY-PT generating series of knot K

- $P^{Q_{K}}(\mathbf{x},q)$  motivic generating series of respective quiver  $Q_{K}$
- $x_i = a^{a_i} q^{q_i C_{ii}} x$  is a change of variables

Knots Quivers Knots-quivers correspondence

#### Knots-quivers correspondence



• Knots-quivers (KQ) correspondence is an equality

$$P^{K}(x,a,q) = P^{Q_{K}}(\mathbf{x},q)\Big|_{x_{i}=a^{a_{i}}q^{q_{i}-C_{ii}}x}$$

- $P^{K}(x, a, q)$  HOMFLY-PT generating series of knot K
- P<sup>Q<sub>K</sub></sup>(x,q) motivic generating series of respective quiver Q<sub>K</sub>
   x<sub>i</sub> = a<sup>a<sub>i</sub></sup> q<sup>q<sub>i</sub>-C<sub>ii</sub>x is a change of variables
  </sup>

ヘロト ヘヨト ヘヨト ヘ

Knots Quivers Knots-quivers correspondence

## Knots-quivers correspondence



• Knots-quivers (KQ) correspondence is an equality

$$P^{K}(x,a,q) = P^{Q_{K}}(\mathbf{x},q)\Big|_{x_{i}=a^{a_{i}}q^{q_{i}-C_{ii}}x}$$

- $P^{K}(x, a, q)$  HOMFLY-PT generating series of knot K
- $P^{Q_{\kappa}}(\mathbf{x},q)$  motivic generating series of respective quiver  $Q_{\kappa}$
- $x_i = a^{a_i} q^{q_i C_{ii}} x$  is a change of variables

・ ロ ト ・ 一戸 ト ・ ・ 三 ト ・ ・

Knots Quivers Knots-quivers correspondence

# Knots-quivers correspondence



• Knots-quivers (KQ) correspondence is an equality

$$P^{K}(x,a,q) = P^{Q_{K}}(\mathbf{x},q)\Big|_{x_{i}=a^{a_{i}}q^{q_{i}-C_{ii}}x}$$

- $P^{K}(x, a, q)$  HOMFLY-PT generating series of knot K
- $P^{Q_{\kappa}}(\mathbf{x},q)$  motivic generating series of respective quiver  $Q_{\kappa}$
- $x_i = a^{a_i} q^{q_i C_{ii}} x$  is a change of variables

Knots Quivers Knots-quivers correspondence

#### KQ correspondence - unknot example



$$P^{0_1}(x, a, q) = \sum_{r=0}^{\infty} \frac{(a^2; q^2)_r}{(q^2; q^2)_r} x^r$$
$$= 1 + \frac{1 - a^2}{1 - q^2} x + \dots$$

$$P^{Q_{0_1}}(x_1, x_2, q) = \sum_{d_1, d_2 \ge 0} (-q)^{d_1^2} \frac{x_1^{d_1}}{(q^2; q^2)_{d_1}} \frac{x_2^{d_2}}{(q^2; q^2)_{d_2}}$$

$$= 1 + \frac{-qx_1 + x_2}{1 - q^2} + \dots$$

イロト イポト イヨト イヨト

900

Knots Quivers Knots-quivers correspondence

#### KQ correspondence - unknot example



• We can check that

$$P^{0_1}(x, a, q) = P^{Q_{0_1}}(x_1, x_2, q) \Big|_{x_1 = a^2 q^{-1} x, x_2 = x}$$
  
$$1 + \frac{1 - a^2}{1 - q^2} x + \ldots = 1 + \frac{-qx_1 + x_2}{1 - q^2} + \ldots \Big|_{x_1 = a^2 q^{-1} x, x_2 = x}$$

• Comparing with  $x_i = a^{a_i} q^{q_i - C_{ii}} x$  we get

 $(a_1, a_2) = (2, 0)$   $(q_1, q_2) = (0, 0)$ 

nar

Knots Quivers Knots-quivers correspondence

#### KQ correspondence - unknot example



• We can check that

$$P^{0_1}(x, a, q) = P^{Q_{0_1}}(x_1, x_2, q) \Big|_{x_1 = a^2 q^{-1} x, x_2 = x}$$
  
$$1 + \frac{1 - a^2}{1 - q^2} x + \ldots = 1 + \frac{-qx_1 + x_2}{1 - q^2} + \ldots \Big|_{x_1 = a^2 q^{-1} x, x_2 = x}$$

• Comparing with  $x_i = a^{a_i} q^{q_i - C_{ii}} x$  we get

$$(a_1, a_2) = (2, 0)$$
  $(q_1, q_2) = (0, 0)$ 

Knots Quivers Knots-quivers correspondence

#### KQ correspondence and BPS states



• We can also look at the level of BPS states:

$$\operatorname{Exp}\left(\frac{N^{K}(x,a,q)}{1-q^{2}}\right) = P^{K} = P^{Q_{K}} = \operatorname{Exp}\left(\frac{\Omega^{Q_{K}}(\mathbf{x},q)}{1-q^{2}}\right)\Big|_{x_{i}=a^{a_{i}}q^{q_{i}-C_{ii}}x}$$
$$N^{K}(x,a,q) = \Omega^{Q_{K}}(\mathbf{x},q)\Big|_{x_{i}=a^{a_{i}}q^{q_{i}-C_{ii}}x}$$

 Integrality of DT invariants for symmetric quivers [Kontsevich-Soibelman, Efimov] implies LMOV conjecture

イロト イポト イヨト イヨト

MQ (P

Knots Quivers Knots-quivers correspondence

# KQ correspondence and BPS states



• We can also look at the level of BPS states:

$$\operatorname{Exp}\left(\frac{N^{K}(x,a,q)}{1-q^{2}}\right) = P^{K} = P^{Q_{K}} = \operatorname{Exp}\left(\frac{\Omega^{Q_{K}}(\mathbf{x},q)}{1-q^{2}}\right)\Big|_{x_{i}=a^{a_{i}}q^{q_{i}-C_{ii}}x}$$
$$N^{K}(x,a,q) = \Omega^{Q_{K}}(\mathbf{x},q)\Big|_{x_{i}=a^{a_{i}}q^{q_{i}-C_{ii}}x}$$

 Integrality of DT invariants for symmetric quivers [Kontsevich-Soibelman, Efimov] implies LMOV conjecture

・ 同 ト ・ ヨ ト ・ ヨ ト

MQ (P

#### General idea Physics Geometry

# Outline



- Knots
- Quivers
- Knots-quivers correspondence

#### 2 Physics and geometry of KQ correspondence

- General idea
- Physics
- Geometry

#### 3 Summary

< □ > < 同 > < 回 > .

- E - N

General idea Physics Geometry

#### KQ correspondence discussed so far





nar



General idea Physics Geometry

#### Towards physics of KQ correspondence





A (10) < (10) </p>

nar



General idea Physics Geometry

# Towards physics of KQ correspondence





イロト イポト イヨト イヨト

900

	Knots		Quivers
Math	HOMFLY-PT gen. series	$\rightarrow$	Motivic gen. series
	$\downarrow$		
Physics	$3d \mathcal{N} = 2 T[L_K]$		

General idea Physics Geometry

#### Missing element





5900



P. Kucharski Physics and geometry of KQ correspondence

(日) (同) (三) (

-
General idea Physics Geometry

#### Missing element





5900



P. Kucharski Physics and geometry of KQ correspondence

(日) (同) (三) (

General idea Physics Geometry

### Missing element





イロト イポト イヨト イヨト

900

Э

	Knots		Quivers
Math	HOMFLY-PT gen. series	$\rightarrow$	Motivic gen. series
	$\downarrow$		$\downarrow$
Physics	3d $\mathcal{N} = 2 T[L_K]$	$\rightarrow$	$3d \mathcal{N} = 2 \mathcal{T}[Q_{\mathcal{K}}]$

General idea Physics Geometry

#### Geometric interpretations



General idea Physics Geometry

## Outline



- Knots
- Quivers
- Knots-quivers correspondence

#### 2 Physics and geometry of KQ correspondence

General idea

#### Physics

Geometry

#### 3 Summary

< ロト < 同ト < ヨト

-

General idea Physics Geometry

### Construction of knot complement theory

• We can construct 3d  $\mathcal{N} = 2$  knot complement theory  $T[M_K]$  basing on large colour and classical limit of HOMFLY-PT polynomial [Fuji, Gukov, Sułkowski]

$$P_{r}^{K}(a,q) \xrightarrow[q^{2r} \text{ fixed}]{\hbar \to 0} \exp\left[\frac{1}{2\hbar} \left(\widetilde{\mathscr{W}}_{\mathcal{T}[M_{K}]} + O(\hbar)\right)\right]$$

•  $\widetilde{\mathscr{W}}_{T[M_{\mathcal{K}}]}$  is a twisted superpotential

$$\operatorname{Li}_2(\ldots) \longleftrightarrow$$
 chiral field  
 $\frac{\kappa}{2} \log(\ldots) \log(\ldots) \longleftrightarrow$  CS coupling



General idea Physics Geometry

# Recall: what is $T[L_K]$

- $T[L_K]$ : 3d  $\mathcal{N} = 2$  eff. theory on  $\mathbb{R}^{1,2}$
- Partition function=generating function of Labastida-Marino-Ooguri-Vafa invariants N<sup>K</sup><sub>r,i,j</sub>

$$Z = \operatorname{Exp}\left(\frac{\sum_{r,i,j} N_{r,i,j}^{K} x^{r} a^{i} q^{j}}{1 - q^{2}}\right)$$

•  $\frac{M5 \text{ brane}}{M2 \text{ branes}} = \frac{\mathbb{R}^{1,2}}{\mathbb{BPS} \text{ states}}$ 

• BPS states are counted by LMOV invariants  $\Rightarrow$  LMOV integrality conjecture



・ 同 ト ・ ヨ ト ・ ヨ

General idea Physics Geometry

# Recall: what is $T[L_K]$

- $T[L_K]$ : 3d  $\mathcal{N} = 2$  eff. theory on  $\mathbb{R}^{1,2}$
- Partition function=generating function of Labastida-Marino-Ooguri-Vafa invariants N<sup>K</sup><sub>r,i,j</sub>

$$Z = \operatorname{Exp}\left(\frac{\sum_{r,i,j} N_{r,i,j}^{K} x^{r} a^{i} q^{j}}{1 - q^{2}}\right)$$

•  $\frac{M5 \text{ brane}}{M2 \text{ branes}} = \frac{\mathbb{R}^{1,2}}{\mathbb{BPS} \text{ states}}$ 

• BPS states are counted by LMOV invariants  $\Rightarrow$  LMOV integrality conjecture



・ 同 ト ・ ヨ ト ・ ヨ

General idea Physics Geometry

# Construction of $T[L_K]$ theory

#### Idea

$$P^{K}(x,a,q) \xrightarrow[q^{2r} \to y]{} \int dy \exp\left[\frac{1}{2\hbar}\left(\widetilde{\mathscr{W}}_{T[L_{K}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathscr{W}}_{T[L_{K}]} = \widetilde{\mathscr{W}}_{T[M_{K}]} + \log x \log y$$

- Integral  $\int dy$  means that U(1) symmetry corresponding to fugacity y is gauged
- We have the same matter content with just extra CS coupling with background field



General idea Physics Geometry

# Construction of $T[L_K]$ theory

#### Idea

$$P^{K}(x,a,q) \xrightarrow[q^{2r} \to y]{} \int dy \exp\left[\frac{1}{2\hbar} \left(\widetilde{\mathscr{W}}_{T[L_{K}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathscr{W}}_{T[L_{K}]} = \widetilde{\mathscr{W}}_{T[M_{K}]} + \log x \log y$$

- Integral  $\int dy$  means that U(1) symmetry corresponding to fugacity y is gauged
- We have the same matter content with just extra CS coupling with background field



General idea Physics Geometry

# Construction of $T[L_K]$ theory

#### Idea

$$P^{K}(x,a,q) \xrightarrow[q^{2r} \to y]{} \int dy \exp\left[\frac{1}{2\hbar} \left(\widetilde{\mathscr{W}}_{T[L_{K}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathscr{W}}_{T[L_{K}]} = \widetilde{\mathscr{W}}_{T[M_{K}]} + \log x \log y$$

- Integral ∫ dy means that U(1) symmetry corresponding to fugacity y is gauged
- We have the same matter content with just extra CS coupling with background field



General idea Physics Geometry

# Construction of $T[L_K]$ theory

#### Idea

$$P^{K}(x,a,q) \xrightarrow[q^{2r} \to y]{} \int dy \exp\left[\frac{1}{2\hbar} \left(\widetilde{\mathscr{W}}_{T[L_{K}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathscr{W}}_{T[L_{K}]} = \widetilde{\mathscr{W}}_{T[M_{K}]} + \log x \log y$$

- Integral  $\int dy$  means that U(1) symmetry corresponding to fugacity y is gauged
- We have the same matter content with just extra CS coupling with background field



# Example: $T[L_{0_1}]$

• For the unknot we have

$$P^{0_{1}}(x, a, q) \xrightarrow[q^{2r} \to y]{h \to 0} \int dy \exp\left[\frac{1}{2\hbar} \left(\widetilde{\mathscr{W}}_{\mathcal{T}[L_{0_{1}}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathscr{W}}_{\mathcal{T}[L_{0_{1}}]} = \operatorname{Li}_{2}(y) + \operatorname{Li}_{2}\left(y^{-1}a^{-2}\right) + \log x \log y$$



- $T[L_{0_1}]$  is a U(1) gauge theory (fugacity y) with one fundamental and one antifundamental chiral
- Antifundamental chiral is charged under the  $U(1)_a$  global symmetry arising from  $S^2$  in the conifold

General idea

Physics

Geometry

(日) (同) (三) (三)

General idea Physics Geometry

# Example: $T[L_{0_1}]$

• For the unknot we have

$$P^{0_1}(x, a, q) \xrightarrow[q^{2r} \to y]{h \to 0} \int dy \exp\left[\frac{1}{2\hbar} \left(\widetilde{\mathscr{W}}_{T[L_{0_1}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathscr{W}}_{T[L_{0_1}]} = \operatorname{Li}_2(y) + \operatorname{Li}_2\left(y^{-1}a^{-2}\right) + \log x \log y$$

- T[L<sub>01</sub>] is a U(1) gauge theory (fugacity y) with one fundamental and one antifundamental chiral
- Antifundamental chiral is charged under the  $U(1)_a$  global symmetry arising from  $S^2$  in the conifold

(日) (同) (三) (三)

rce euclidials e

General idea Physics Geometry

# Example: $T[L_{0_1}]$

• For the unknot we have

$$P^{0_1}(x, a, q) \xrightarrow[q^{2r} \to y]{h \to 0} \int dy \exp\left[\frac{1}{2\hbar} \left(\widetilde{\mathscr{W}}_{T[L_{0_1}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathscr{W}}_{T[L_{0_1}]} = \operatorname{Li}_2(y) + \operatorname{Li}_2\left(y^{-1}a^{-2}\right) + \log x \log y$$

- $T[L_{0_1}]$  is a U(1) gauge theory (fugacity y) with one fundamental and one antifundamental chiral
- Antifundamental chiral is charged under the  $U(1)_a$  global symmetry arising from  $S^2$  in the conifold

< ロ > ( 同 > ( 回 > ( 回 > ))

rce euclidials e

General idea Physics Geometry

### Recall of general idea - now we are here





イロト イポト イヨト イヨト

5900

	Knots		Quivers
Math	HOMFLY-PT gen. series	$\rightarrow$	Motivic gen. series
	$\downarrow$		
Physics	$3d \mathcal{N} = 2 T[L_K]$		

General idea Physics Geometry

## We look for the missing element





イロト イポト イヨト イヨ

5900

	Knots		Quivers
Math	HOMFLY-PT gen. series	$\rightarrow$	Motivic gen. series
	$\downarrow$		$\downarrow$
Physics	3d $\mathcal{N} = 2 T[L_K]$		

General idea Physics Geometry

## We look for the missing element





イロト イポト イヨト イヨ

5900

	Knots		Quivers
Math	HOMFLY-PT gen. series	$\rightarrow$	Motivic gen. series
	$\downarrow$		$\downarrow$
Physics	3d $\mathcal{N} = 2 T[L_K]$		$3d \mathcal{N} = 2 T[Q_K]$

General idea Physics Geometry

# Missing element: $T[Q_K]$ theory

#### Idea

Consider large colour and classical limit of motivic generating series!

$$P^{Q_{K}}(\mathbf{x},q) \xrightarrow[q^{2d_{i}} \rightarrow y_{i}]{} \int \prod_{i} dy_{i} \exp\left[\frac{1}{2\hbar}\left(\widetilde{\mathcal{W}}_{\mathcal{T}[Q_{K}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathcal{W}}_{\mathcal{T}[Q_{K}]} = \sum_{i} \left[\operatorname{Li}_{2}\left(y_{i}\right) + \log x_{i} \log y_{i}\right] + \sum_{i,j} \frac{C_{i,j}}{2} \log y_{i} \log y_{j}.$$

- Gauge group:  $U(1)^{\text{#vertices}}$
- Matter content: one chiral for each vertex

P Kucharski

• CS couplings given by  $C_{i,j} = #$ arrows

WS

Physics and geometry of KQ correspondence Summary General idea Physics Geometry

# Missing element: $T[Q_K]$ theory

#### Idea

Consider large colour and classical limit of motivic generating series!

$$P^{Q_{\kappa}}(\mathbf{x},q) \xrightarrow[q^{2d_{i}} \rightarrow y_{i}]{} \int \prod_{i} dy_{i} \exp\left[\frac{1}{2\hbar}\left(\widetilde{\mathscr{W}}_{T[Q_{\kappa}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathscr{W}}_{T[Q_{\kappa}]} = \sum_{i} \left[\operatorname{Li}_{2}\left(y_{i}\right) + \log x_{i} \log y_{i}\right] + \sum_{i,j} \frac{C_{i,j}}{2} \log y_{i} \log y_{j}$$

- Gauge group:  $U(1)^{\text{#vertices}}$
- Matter content: one chiral for each vertex

P Kucharski

• CS couplings given by  $C_{i,j} = \#$ arrows

Physics and geometry of KQ correspondence



Physics and geometry of KQ correspondence Summary General idea Physics Geometry

## Missing element: $T[Q_K]$ theory

#### Idea

Consider large colour and classical limit of motivic generating series!

$$P^{Q_{\kappa}}(\mathbf{x},q) \xrightarrow[q^{2d_{i}} \rightarrow y_{i}]{} \int \prod_{i} dy_{i} \exp\left[\frac{1}{2\hbar}\left(\widetilde{\mathscr{W}}_{T[Q_{\kappa}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathscr{W}}_{T[Q_{\kappa}]} = \sum_{i} \left[\operatorname{Li}_{2}\left(y_{i}\right) + \log x_{i} \log y_{i}\right] + \sum_{i,j} \frac{C_{i,j}}{2} \log y_{i} \log y_{j}$$

- Gauge group:  $U(1)^{\text{#vertices}}$
- Matter content: one chiral for each vertex
- CS couplings given by  $C_{i,j} = \#$ arrows

they.

General idea Physics Geometry

## Missing element: $T[Q_K]$ theory

#### Idea

Consider large colour and classical limit of motivic generating series!

$$P^{Q_{\kappa}}(\mathbf{x},q) \xrightarrow[q^{2d_{i}} \rightarrow y_{i}]{} \int \prod_{i} dy_{i} \exp\left[\frac{1}{2\hbar}\left(\widetilde{\mathscr{W}}_{T[Q_{\kappa}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathscr{W}}_{T[Q_{\kappa}]} = \sum_{i} \left[\operatorname{Li}_{2}\left(y_{i}\right) + \log x_{i} \log y_{i}\right] + \sum_{i,j} \frac{C_{i,j}}{2} \log y_{i} \log y_{j},$$

- Gauge group:  $U(1)^{\text{#vertices}}$
- Matter content: one chiral for each vertex

• CS couplings given by  $C_{i,j} = #$ arrows

Physics and geometry of KQ correspondence Summary General idea Physics Geometry

## Missing element: $T[Q_K]$ theory

#### Idea

Consider large colour and classical limit of motivic generating series!

$$P^{Q_{\kappa}}(\mathbf{x},q) \xrightarrow[q^{2d_{i}} \rightarrow y_{i}]{} \int \prod_{i} dy_{i} \exp\left[\frac{1}{2\hbar}\left(\widetilde{\mathscr{W}}_{T[Q_{\kappa}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathscr{W}}_{T[Q_{\kappa}]} = \sum_{i} \left[\operatorname{Li}_{2}\left(y_{i}\right) + \log x_{i} \log y_{i}\right] + \sum_{i,j} \frac{C_{i,j}}{2} \log y_{i} \log y_{j}$$

- Gauge group:  $U(1)^{\text{#vertices}}$
- Matter content: one chiral for each vertex

• CS couplings given by 
$$C_{i,j} = \#$$
arrows

and a

General idea Physics Geometry

# Example: $T[Q_{0_1}]$

• For the unknot quiver we have

$$P^{Q_{0_{1}}}(x_{1}, x_{2}, q) \xrightarrow[q^{2d_{i}} \to y_{i}]{} \int dy_{1} dy_{2} \exp\left[\frac{1}{2\hbar}\left(\widetilde{\mathscr{W}}_{\mathcal{T}[Q_{0_{1}}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathscr{W}}_{\mathcal{T}[Q_{0_{1}}]} = \operatorname{Li}_{2}(y_{1}) + \operatorname{Li}_{2}(y_{2}) + \log x_{1} \log y_{1} + \log x_{2} \log y_{2} + \frac{1}{2} \log y_{1} \log y_{1}$$

\$\mathcal{T}[Q\_{0\_1}]\$ is a \$U(1)^{(1)} \times U(1)^{(2)}\$ gauge theory with one chiral field for each group

2

• CS level one for  $U(1)^{(1)}$ , consistent with  $C^{0_1} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$ 

General idea Physics Geometry

# Example: $T[Q_{0_1}]$

• For the unknot quiver we have

$$P^{Q_{0_{1}}}(x_{1}, x_{2}, q) \xrightarrow[q^{2d_{i}} \to y_{i}]{} \int dy_{1} dy_{2} \exp\left[\frac{1}{2\hbar}\left(\widetilde{\mathscr{W}}_{T[Q_{0_{1}}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathscr{W}}_{T[Q_{0_{1}}]} = \operatorname{Li}_{2}(y_{1}) + \operatorname{Li}_{2}(y_{2}) + \log x_{1} \log y_{1} + \log x_{2} \log y_{2} + \frac{1}{2} \log y_{1} \log y_{1}$$

•  $\mathcal{T}[Q_{0_1}]$  is a  $U(1)^{(1)} \times U(1)^{(2)}$  gauge theory with one chiral field for each group

• CS level one for  $U(1)^{(1)}$ , consistent with  $C^{0_1} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$ 

(日) (同) (三) (

General idea Physics Geometry

# Example: $T[Q_{0_1}]$

• For the unknot quiver we have

$$P^{Q_{0_{1}}}(x_{1}, x_{2}, q) \xrightarrow[q^{2d_{i}} \to y_{i}]{} \int dy_{1} dy_{2} \exp\left[\frac{1}{2\hbar}\left(\widetilde{\mathscr{W}}_{T[Q_{0_{1}}]} + O(\hbar)\right)\right]$$
$$\widetilde{\mathscr{W}}_{T[Q_{0_{1}}]} = \operatorname{Li}_{2}(y_{1}) + \operatorname{Li}_{2}(y_{2}) + \log x_{1} \log y_{1} + \log x_{2} \log y_{2} + \frac{1}{2} \log y_{1} \log y_{1}$$

•  $\mathcal{T}[Q_{0_1}]$  is a  $U(1)^{(1)} \times U(1)^{(2)}$  gauge theory with one chiral field for each group



• CS level one for  $U(1)^{(1)}$ , consistent with  $C^{0_1} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$ 

Image: A matrix

Physics and geometry of KQ correspondence Summary General idea Physics Geometry

### Recall of general idea - now we are here





イロト イポト イヨト イヨト

900

	Knots		Quivers
Math	HOMFLY-PT gen. series	$\rightarrow$	Motivic gen. series
	$\downarrow$		$\downarrow$
Physics	3d $\mathcal{N} = 2 T[L_K]$		$3d \mathcal{N} = 2 \mathcal{T}[Q_{\mathcal{K}}]$

General idea Physics Geometry

We want the physical interpretation of KQ correspondence





	Knots		Quivers
Math	HOMFLY-PT gen. series	$\rightarrow$	Motivic gen. series
	$\downarrow$		$\downarrow$
Physics	3d $\mathcal{N} = 2 T[L_K]$	$\rightarrow$	$3d \mathcal{N} = 2 \mathcal{T}[Q_{\mathcal{K}}]$

P. Kucharski Physics and geometry of KQ correspondence

(日) (同) (三) (

DQ P

#### General idea Physics Geometry

# Duality

 Physical meaning of the knots-quivers correspondence is a duality between two 3d *N* = 2 theories:

 $T[L_K] \longleftrightarrow T[Q_K]$ 



• We already know that

$$Z(T[L_K]) = P^K(x, a, q)$$

• We can calculate the partition function of  $\mathcal{T}[Q_K]$  and indeed

$$Z(T[Q_K]) = P^{Q_K}(\mathbf{x}, q)$$

#### General idea Physics Geometry

# Duality

 Physical meaning of the knots-quivers correspondence is a duality between two 3d *N* = 2 theories:

 $T[L_K] \longleftrightarrow T[Q_K]$ 



• We already know that

$$Z(T[L_K]) = P^K(x, a, q)$$

• We can calculate the partition function of  $\mathcal{T}[Q_K]$  and indeed

$$Z(T[Q_K]) = P^{Q_K}(\mathbf{x}, q)$$

General idea Physics Geometry

# Duality

 Physical meaning of the knots-quivers correspondence is a duality between two 3d *N* = 2 theories:

 $T[L_K] \longleftrightarrow T[Q_K]$ 



3 b. 4

• We already know that

$$Z(T[L_K]) = P^K(x, a, q)$$

• We can calculate the partition function of  $\mathcal{T}[Q_{\mathcal{K}}]$  and indeed

$$Z(T[Q_{\mathcal{K}}]) = P^{Q_{\mathcal{K}}}(\mathbf{x},q)$$

## Outline



- Knots
- Quivers
- Knots-quivers correspondence

#### 2 Physics and geometry of KQ correspondence

- General idea
- Physics
- Geometry

#### 3 Summary

< ロト < 同ト < 三ト

- E - N

Quiver nodes extracted from motivic generating series

• Let's look at motivic generating series restricted to  $|\mathbf{d}| = 1$ 

$$P^{Q_{K}}_{|\mathbf{d}|=1}(\mathsf{x},q) = \sum_{i \in Q_{0}} rac{(-q)^{C_{ii}} x_{i}}{1-q^{2}}$$

- Every vertex  $(i \in Q_0)$  contributes once
- No interactions between different nodes
- What is the meaning of  $P_{|\mathbf{d}|=1}^{Q_{K}}(\mathbf{x},q)$ ?



(日) (同) (三) (三) (三)

Quiver nodes extracted from motivic generating series

• Let's look at motivic generating series restricted to  $|\mathbf{d}| = 1$ 

$$P_{|\mathbf{d}|=1}^{Q_{\mathcal{K}}}(\mathsf{x},q) = \sum_{i \in Q_0} rac{(-q)^{\mathcal{C}_{ii}} x_i}{1-q^2}$$

- Every vertex  $(i \in Q_0)$  contributes once
- No interactions between different nodes
- What is the meaning of  $P_{|\mathbf{d}|=1}^{Q_K}(\mathbf{x},q)$ ?

イロト イポト イラト イラト

Quiver nodes extracted from motivic generating series

• Let's look at motivic generating series restricted to  $|\mathbf{d}| = 1$ 

$$P^{Q_{K}}_{|\mathbf{d}|=1}(\mathsf{x},q) = \sum_{i \in Q_{0}} rac{(-q)^{C_{ii}} x_{i}}{1-q^{2}}$$

- Every vertex  $(i \in Q_0)$  contributes once
- No interactions between different nodes
- What is the meaning of  $P_{|\mathbf{d}|=1}^{Q_K}(\mathbf{x},q)$ ?



イロト イポト イヨト イヨト

Quiver nodes extracted from motivic generating series

• Let's look at motivic generating series restricted to  $|\mathbf{d}| = 1$ 

$$P_{|\mathbf{d}|=1}^{Q_{\mathcal{K}}}(\mathsf{x},q) = \sum_{i \in Q_0} rac{(-q)^{\mathcal{C}_{ii}} x_i}{1-q^2}$$

- Every vertex  $(i \in Q_0)$  contributes once
- No interactions between different nodes
- What is the meaning of  $P_{|\mathbf{d}|=1}^{Q_{K}}(\mathbf{x},q)$ ?

イロト イポト イラト イラト

General idea Physics Geometry

#### Quiver nodes as homology generators

• Let's apply KQ change of variables

$$P_{|\mathbf{d}|=1}^{Q_{K}}(\mathbf{x},q)\Big|_{x_{i}=a^{a_{i}}q^{q_{i}-c_{iix}}}=P_{1}^{K}(a,q)x$$

*P*<sup>K</sup><sub>1</sub>(*a*, *q*) is an Euler characteristic of HOMFLY-PT homology *H*(*K*) with set of generators *G*(*K*), so

$$\sum_{\substack{\in Q_0}} \frac{(-q)^{C_{ii}} x_i}{1-q^2} \bigg|_{x_i = a^{a_i} q^{q_i - C_{ii}} x} = \frac{\sum_{i \in \mathscr{G}(K)} a^{a_i} q^{q_i} (-1)^{C_{ii}}}{1-q^2} x$$



• KQ change of variables is encoded in homological degrees
General idea Physics Geometry

#### Quiver nodes as homology generators

• Let's apply KQ change of variables

$$P_{|\mathbf{d}|=1}^{Q_{K}}(\mathbf{x},q)\Big|_{x_{i}=a^{a_{i}}q^{q_{i}-C_{ii}}\times}=P_{1}^{K}(a,q)x$$

*P*<sup>K</sup><sub>1</sub>(*a*, *q*) is an Euler characteristic of HOMFLY-PT homology *H*(*K*) with set of generators *G*(*K*), so

$$\sum_{i \in Q_0} \frac{(-q)^{C_{ii}} x_i}{1-q^2} \bigg|_{x_i = a^{a_i} q^{q_i - C_{ii}} x} = \frac{\sum_{i \in \mathscr{G}(K)} a^{a_i} q^{q_i} (-1)^{C_{ii}}}{1-q^2} x$$

• Each vertex corresponds to the homology generator

• KQ change of variables is encoded in homological degrees

P. Kucharski Physics and geometry of KQ correspondence

General idea Physics Geometry

#### Quiver nodes as homology generators

• Let's apply KQ change of variables

$$P_{|\mathbf{d}|=1}^{Q_{K}}(\mathbf{x},q)\Big|_{x_{i}=a^{a_{i}}q^{q_{i}-C_{ii}}\times}=P_{1}^{K}(a,q)x$$

*P*<sup>K</sup><sub>1</sub>(*a*, *q*) is an Euler characteristic of HOMFLY-PT homology *H*(*K*) with set of generators *G*(*K*), so

$$\sum_{i \in Q_0} \frac{(-q)^{C_{ii}} x_i}{1-q^2} \bigg|_{x_i = a^{a_i} q^{q_i - C_{ii}} x} = \frac{\sum_{i \in \mathscr{G}(\kappa)} a^{a_i} q^{q_i} (-1)^{C_{ii}}}{1-q^2} x$$



• KQ change of variables is encoded in homological degrees



General idea Physics Geometry

#### Quiver nodes as homology generators

• Let's apply KQ change of variables

$$P_{|\mathbf{d}|=1}^{Q_{K}}(\mathbf{x},q)\Big|_{x_{i}=a^{a_{i}}q^{q_{i}-c_{ii}}}=P_{1}^{K}(a,q)x$$

*P*<sup>K</sup><sub>1</sub>(*a*, *q*) is an Euler characteristic of HOMFLY-PT homology *H*(*K*) with set of generators *G*(*K*), so

$$\sum_{i \in Q_0} \frac{(-q)^{C_{ii}} x_i}{1-q^2} \bigg|_{x_i = a^{a_i} q^{q_i - C_{ii}} x} = \frac{\sum_{i \in \mathscr{G}(K)} a^{a_i} q^{q_i} (-1)^{C_{ii}}}{1-q^2} x$$



• KQ change of variables is encoded in homological degrees



General idea Physics Geometry

### Quiver nodes as BPS states

• Generators of uncoloured homology correspond to BPS states [Gukov, Schwarz, Vafa]



• Quiver nodes correspond to BPS states of  $\mathcal{T}[L_{\mathcal{K}}]$  counted by LMOV invariants



• ...or BPS states of  $T[Q_K]$  counted by DT invariants





イロト イポト イヨト イヨ

General idea Physics Geometry

### Quiver nodes as BPS states

• Generators of uncoloured homology correspond to BPS states [Gukov, Schwarz, Vafa]

• Quiver nodes correspond to BPS states of  $T[L_K]$  counted by LMOV invariants

$$\sum_{i \in Q_0} \frac{(-q)^{C_{ii}} x_i}{1-q^2} \bigg|_{x_i=a^{a_i}q^{q_i-C_{ii}} x} = \frac{N_1^K(a,q)}{1-q^2} x$$

• ... or BPS states of  $T[Q_K]$  counted by DT invariants

$$\sum_{i \in Q_0} \frac{(-q)^{C_{ii}}_{X_i}}{1-q^2} = \frac{\Omega_{|\mathbf{d}|=1}^{Q_K}(\mathbf{x},q)}{1-q^2}$$



General idea Physics Geometry

### Quiver nodes as BPS states

• Generators of uncoloured homology correspond to BPS states [Gukov, Schwarz, Vafa]

• Quiver nodes correspond to BPS states of  $T[L_K]$  counted by LMOV invariants

$$\sum_{i \in Q_0} \frac{(-q)^{C_{ii}} x_i}{1-q^2} \bigg|_{x_i=a^{a_i}q^{q_i-C_{ii}} \times} = \frac{N_1^K(a,q)}{1-q^2} x$$

• ...or BPS states of  $T[Q_K]$  counted by DT invariants

$$\sum_{i \in Q_0} \frac{(-q)^{C_{ii}} x_i}{1-q^2} = \frac{\Omega_{|\mathbf{d}|=1}^{Q_{\mathcal{K}}}(\mathbf{x},q)}{1-q^2}$$

A (1) < (1) < (1) < (1) < (1) </p>



General idea Physics Geometry

### Quiver nodes as holomorphic disks

- Recall: M2 branes are BPS states on the flat side and holomorphic disks on the Calabi-Yau side
- We can reinterpret topological data of KQ change of variables in the language of disks:
  - power r in  $x^r \longrightarrow \#$ windings around  $L_K$
  - power  $a_i$  in  $a^{a_i} \longrightarrow \#$ wrappings around base  $S^2$
  - power  $q_i$  in  $q^{q_i} \longrightarrow$  invariant self-linking#
  - $C_{ii} = t_i \longrightarrow \text{linking} \#$  between disk and its small shift
- Quiver nodes correspond to basic disks holomorphic curves that wind around *L<sub>K</sub>* once (*r* = 1)

< ロ > < 同 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ >



General idea Physics Geometry

# Quiver nodes as holomorphic disks

- Recall: M2 branes are BPS states on the flat side and holomorphic disks on the Calabi-Yau side
- We can reinterpret topological data of KQ change of variables in the language of disks:
  - power r in  $x^r \longrightarrow \#$ windings around  $L_K$
  - power  $a_i$  in  $a^{a_i} \longrightarrow \#$ wrappings around base  $S^2$
  - power  $q_i$  in  $q^{q_i} \longrightarrow$  invariant self-linking#
  - $C_{ii} = t_i \longrightarrow \text{linking} \#$  between disk and its small shift
- Quiver nodes correspond to basic disks holomorphic curves that wind around *L<sub>K</sub>* once (*r* = 1)

< ロ > < 同 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ >



General idea Physics Geometry

# Quiver nodes as holomorphic disks

- Recall: M2 branes are BPS states on the flat side and holomorphic disks on the Calabi-Yau side
- We can reinterpret topological data of KQ change of variables in the language of disks:
  - power r in  $x^r \longrightarrow \#$ windings around  $L_K$
  - power  $a_i$  in  $a^{a_i} \longrightarrow \#$ wrappings around base  $S^2$
  - power  $q_i$  in  $q^{q_i} \longrightarrow$  invariant self-linking#
  - $C_{ii} = t_i \longrightarrow \text{linking} \#$  between disk and its small shift
- Quiver nodes correspond to basic disks holomorphic curves that wind around *L<sub>K</sub>* once (*r* = 1)

< ロ > < 同 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ >



General idea Physics Geometry

# Quiver nodes as holomorphic disks

- Recall: M2 branes are BPS states on the flat side and holomorphic disks on the Calabi-Yau side
- We can reinterpret topological data of KQ change of variables in the language of disks:
  - power r in  $x^r \longrightarrow \#$ windings around  $L_K$
  - power  $a_i$  in  $a^{a_i} \longrightarrow \#$ wrappings around base  $S^2$
  - power  $q_i$  in  $q^{q_i} \longrightarrow$  invariant self-linking#
  - $C_{ii} = t_i \longrightarrow \text{linking} \#$  between disk and its small shift
- Quiver nodes correspond to basic disks holomorphic curves that wind around *L<sub>K</sub>* once (*r* = 1)



General idea Physics Geometry

# Quiver nodes as holomorphic disks

- Recall: M2 branes are BPS states on the flat side and holomorphic disks on the Calabi-Yau side
- We can reinterpret topological data of KQ change of variables in the language of disks:
  - power r in  $x^r \longrightarrow \#$ windings around  $L_K$
  - power  $a_i$  in  $a^{a_i} \longrightarrow \#$ wrappings around base  $S^2$
  - power  $q_i$  in  $q^{q_i} \longrightarrow$  invariant self-linking#
  - $C_{ii} = t_i \longrightarrow \text{linking} \#$  between disk and its small shift
- Quiver nodes correspond to basic disks holomorphic curves that wind around *L<sub>K</sub>* once (*r* = 1)



< 日 > < 同 > < 三 > < 三 > < 三 > <

General idea Physics Geometry

# Quiver nodes as holomorphic disks

- Recall: M2 branes are BPS states on the flat side and holomorphic disks on the Calabi-Yau side
- We can reinterpret topological data of KQ change of variables in the language of disks:
  - power r in  $x^r \longrightarrow \#$ windings around  $L_K$
  - power  $a_i$  in  $a^{a_i} \longrightarrow \#$ wrappings around base  $S^2$
  - power  $q_i$  in  $q^{q_i} \longrightarrow$  invariant self-linking#
  - $C_{ii} = t_i \longrightarrow \text{linking} \#$  between disk and its small shift
- Quiver nodes correspond to basic disks holomorphic curves that wind around *L<sub>K</sub>* once (*r* = 1)



・ロト ・ 一下・ ・ 日 ・ ・ 日 ・

General idea Physics Geometry

# Quiver nodes as holomorphic disks

- Recall: M2 branes are BPS states on the flat side and holomorphic disks on the Calabi-Yau side
- We can reinterpret topological data of KQ change of variables in the language of disks:
  - power r in  $x^r \longrightarrow \#$ windings around  $L_K$
  - power  $a_i$  in  $a^{a_i} \longrightarrow \#$ wrappings around base  $S^2$
  - power  $q_i$  in  $q^{q_i} \longrightarrow$  invariant self-linking#
  - $C_{ii} = t_i \longrightarrow \text{linking} \#$  between disk and its small shift
- Quiver nodes correspond to basic disks holomorphic curves that wind around  $L_K$  once (r = 1)



General idea Physics **Geometry** 

### Example: unknot quiver vertices and homology

• For the unknot quiver we have

$$P^{Q_{\mathbf{0}_1}}_{|\mathbf{d}|=1}(\mathbf{x},q) = rac{-qx_1+x_2}{1-q^2}$$

•  $\mathscr{H}(0_1)$  has two generators with degrees

$$(a_1, a_2) = (2, 0), \quad (q_1, q_2) = (0, 0)$$
  
 $(t_1, t_2) = (1, 0) = (C_{11}, C_{22})$ 



 Generators correspond to quiver vertices and their degrees encode KQ change of variables x<sub>1</sub> = a<sup>2</sup>q<sup>-1</sup>x, x<sub>2</sub> = x giving

$$P_1^{0_1}(a,q) \times = \frac{-a^2 x + x}{1 - q^2} = \frac{\sum_{i \in \mathscr{G}(0_1)} a^{a_i} q^{q_i} (-1)^{C_{ii}}}{1 - q^2} \times$$

P. Kucharski Physics and geometry of KQ correspondence

General idea Physics **Geometry** 

Example: unknot quiver vertices and homology

• For the unknot quiver we have

$$P^{Q_{\mathbf{0}_1}}_{|\mathbf{d}|=1}(\mathbf{x},q) = rac{-qx_1+x_2}{1-q^2}$$

•  $\mathscr{H}(0_1)$  has two generators with degrees

$$(a_1, a_2) = (2, 0), \quad (q_1, q_2) = (0, 0)$$
  
 $(t_1, t_2) = (1, 0) = (C_{11}, C_{22})$ 



 Generators correspond to quiver vertices and their degrees encode KQ change of variables x<sub>1</sub> = a<sup>2</sup>q<sup>-1</sup>x, x<sub>2</sub> = x giving

$$P_1^{0_1}(a,q)x = \frac{-a^2x + x}{1 - q^2} = \frac{\sum_{i \in \mathscr{G}(0_1)} a^{a_i} q^{q_i} (-1)^{C_{ii}}}{1 - q^2} x$$

P. Kucharski Physics and geometry of KQ correspondence

Physics and geometry of KQ correspondence Summary

General idea Physics Geometry

Example: unknot quiver vertices and homology

• For the unknot quiver we have

$$P^{Q_{\mathbf{0}_1}}_{|\mathbf{d}|=1}(\mathbf{x},q) = rac{-qx_1+x_2}{1-q^2}$$

•  $\mathcal{H}(0_1)$  has two generators with degrees

$$(a_1, a_2) = (2, 0), \quad (q_1, q_2) = (0, 0)$$
  
 $(t_1, t_2) = (1, 0) = (C_{11}, C_{22})$ 



• Generators correspond to quiver vertices and their degrees encode KQ change of variables  $x_1 = a^2 q^{-1} x$ ,  $x_2 = x$  giving

P. Kucharski

$$P_1^{0_1}(a,q)x = \frac{-a^2x + x}{1-q^2} = \frac{\sum_{i \in \mathscr{G}(0_1)} a^{a_i} q^{q_i} (-1)^{C_{ii}}}{1-q^2} x$$
  
P. Kucharski
Physics and geometry of KQ correspondence

### Example: unknot quiver vertices and BPS states

- Two generators of  $\mathcal{H}(0_1)$  correspond to two BPS states
- BPS states of *T*[*L*<sub>01</sub>] are counted by LMOV invariants

$$N_1^{0_1}(a,q) = 1 - a^2$$

• BPS states of  $T[Q_{0_1}]$  are counted by DT invariants

$$\Omega_{|\mathbf{d}|=1}^{Q_{0_1}}(\mathbf{x},q) = -qx_1 + x_2$$



イロト イポト イラト イラ

### Example: unknot quiver vertices and BPS states

- Two generators of  $\mathcal{H}(0_1)$  correspond to two BPS states
- BPS states of  $T[L_{0_1}]$  are counted by LMOV invariants

$$N_1^{0_1}(a,q) = 1 - a^2$$

• BPS states of  $T[Q_{0_1}]$  are counted by DT invariants

$$\Omega_{|\mathbf{d}|=1}^{Q_{0_1}}(\mathbf{x},q) = -qx_1 + x_2$$



イロト イポト イラト イラ

Summary Geometry

### Example: unknot quiver vertices and BPS states

- Two generators of  $\mathscr{H}(0_1)$  correspond to two BPS states
- BPS states of  $T[L_{0_1}]$  are counted by LMOV invariants

$$N_1^{0_1}(a,q) = 1 - a^2$$

• BPS states of  $T[Q_{0_1}]$  are counted by DT invariants

$$\Omega^{Q_{0_1}}_{|\mathbf{d}|=1}(\mathbf{x},q) = -qx_1 + x_2$$



(1日) (1日) (1日)

Example: unknot quiver vertices and holomorphic disks

- $\bullet\,$  Two BPS states on the flat side  $\longleftrightarrow$  two holomorphic disks on the Calabi-Yau side
- For the first we have the following intepretation of topological data in KQ change of variables:
- $r = 1 \longrightarrow$  winding around  $L_K$
- $a_1 = 2 \longrightarrow$  wrappings around base  $S^2$
- $q_1 = 0 \longrightarrow$  invariant self-linking
- $C_{11} = 1 \longrightarrow$  linking between disk and its small shift



< ロ > ( 同 > ( 回 > ( 回 > ))

Example: unknot quiver vertices and holomorphic disks

- $\bullet\,$  Two BPS states on the flat side  $\longleftrightarrow$  two holomorphic disks on the Calabi-Yau side
- For the first we have the following intepretation of topological data in KQ change of variables:
- $r = 1 \longrightarrow$  winding around  $L_K$
- $a_1 = 2 \longrightarrow$  wrappings around base  $S^2$
- $q_1 = 0 \longrightarrow$  invariant self-linking
- $C_{11} = 1 \longrightarrow$  linking between disk and its small shift



化口石 化晶石 化晶石 化晶石

Example: unknot quiver vertices and holomorphic disks

- $\bullet\,$  Two BPS states on the flat side  $\longleftrightarrow$  two holomorphic disks on the Calabi-Yau side
- For the second we have the following intepretation of topological data in KQ change of variables:
- $r = 1 \longrightarrow$  winding around  $L_K$
- $a_2 = 0 \longrightarrow$  wrappings around base  $S^2$
- $q_2 = 0 \longrightarrow$  invariant self-linking
- $C_{22} = 0 \longrightarrow$  linking between disk and its small shift



< ロ > < 同 > < 回 > < 回 > < 回 > <

Quiver arrows extracted from motivic generating series

$${\mathcal P}^Q({\sf x},q) := \sum_{d_1,...,d_{|\mathcal{Q}_0|} \ge 0} (-q)^{\sum_{1 \le i,j \le |\mathcal{Q}_0|} {\mathcal C}_{i,j} d_i d_j} \prod_{i=1}^{|\mathcal{Q}_0|} rac{x_i^{d_i}}{(q^2;q^2)_{d_i}}$$

- Dimension *d<sub>i</sub>* encodes the number of factors corresponding to *i*-th vertex
- #arrows C<sub>i,j</sub> encodes interactions between vertices i & j
- Motivic generating series counts all objects that can be made from basic ones (nodes) according to quiver arrows



Quiver arrows extracted from motivic generating series

$$P^Q(\mathsf{x},q) := \sum_{d_1,...,d_{|\mathcal{Q}_0|} \ge 0} (-q)^{\sum_{1 \le i,j \le |\mathcal{Q}_0|} \mathcal{C}_{i,j} d_i d_j} \prod_{i=1}^{|\mathcal{Q}_0|} rac{x_i^{d_i}}{(q^2;q^2)_{d_i}}$$

- Dimension *d<sub>i</sub>* encodes the number of factors corresponding to *i*-th vertex
- #arrows C<sub>i,j</sub> encodes interactions between vertices i & j
- Motivic generating series counts all objects that can be made from basic ones (nodes) according to quiver arrows



Quiver arrows extracted from motivic generating series

$${\mathcal P}^Q({\sf x},q) := \sum_{d_1,...,d_{|{\mathcal Q}_0|} \ge 0} (-q)^{\sum_{1 \le i,j \le |{\mathcal Q}_0|} {\mathcal C}_{i,j} d_i d_j} \prod_{i=1}^{|{\mathcal Q}_0|} rac{x_i^{d_i}}{(q^2;q^2)_{d_i}}$$

- Dimension *d<sub>i</sub>* encodes the number of factors corresponding to *i*-th vertex
- #arrows C<sub>i,j</sub> encodes interactions between vertices i & j
- Motivic generating series counts all objects that can be made from basic ones (nodes) according to quiver arrows



Quiver arrows extracted from motivic generating series

$${\mathcal P}^Q({\sf x},q) := \sum_{d_1,...,d_{|{\mathcal Q}_0|} \ge 0} (-q)^{\sum_{1 \le i,j \le |{\mathcal Q}_0|} {\mathcal C}_{i,j} d_i d_j} \prod_{i=1}^{|{\mathcal Q}_0|} rac{x_i^{d_i}}{(q^2;q^2)_{d_i}}$$

- Dimension *d<sub>i</sub>* encodes the number of factors corresponding to *i*-th vertex
- #arrows C<sub>i,j</sub> encodes interactions between vertices i & j
- Motivic generating series counts all objects that can be made from basic ones (nodes) according to quiver arrows



General idea Physics Geometry

### Quiver arrows as disk intersections





- Geometrically P<sup>Q</sup>(x, q) counts all holomorphic curves that can be made from basic disks according to quiver arrows
- Dimension vector  $d_i \longrightarrow \#$ copies of the disk
- #arrows  $C_{i,j} \longrightarrow$  disk boundaries linking#
- Example: one pair of arrows corresponds to two bagel disk boundaries with linking# = 1

イロト イポト イヨト イヨト

General idea Physics Geometry

### Quiver arrows as disk intersections





- Geometrically P<sup>Q</sup>(x, q) counts all holomorphic curves that can be made from basic disks according to quiver arrows
- Dimension vector  $d_i \longrightarrow \#$ copies of the disk
- #arrows  $C_{i,j} \longrightarrow$  disk boundaries linking#
- Example: one pair of arrows corresponds to two bagel disk boundaries with linking# = 1

A B > A B > A B >
 A
 B >
 A
 B >
 A
 B >
 A
 B >
 A
 B >
 A
 B >
 A
 B >
 A
 B >
 A
 B >
 A
 B >
 A
 B >
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A
 B
 A

General idea Physics Geometry

### Quiver arrows as disk intersections





- Geometrically P<sup>Q</sup>(x, q) counts all holomorphic curves that can be made from basic disks according to quiver arrows
- Dimension vector  $d_i \longrightarrow \#$  copies of the disk
- #arrows  $C_{i,j} \longrightarrow$  disk boundaries linking#
- Example: one pair of arrows corresponds to two bagel disk boundaries with linking# = 1

・ロト ・ 日 ・ ・ 日 ・ ・ 日 ・

General idea Physics Geometry

### Quiver arrows as disk intersections





- Geometrically P<sup>Q</sup>(x, q) counts all holomorphic curves that can be made from basic disks according to quiver arrows
- Dimension vector  $d_i \longrightarrow \#$  copies of the disk
- #arrows  $C_{i,j} \longrightarrow$  disk boundaries linking#
- Example: one pair of arrows corresponds to two bagel disk boundaries with linking# = 1

・ロット 全部 マート・ キャー

- Physically knots-quivers correspondence is a duality between 3d  $\mathcal{N} = 2$  theories  $\mathcal{T}[L_K]$  and  $\mathcal{T}[Q_K]$
- Quiver elements can be intepreted in terms of  $T[Q_K]$  data as well as holomorphic disks:



• More details in the paper - coming soon!

・ロト ・ 同ト ・ ヨト ・ ヨト

- Physically knots-quivers correspondence is a duality between 3d  $\mathcal{N} = 2$  theories  $\mathcal{T}[L_K]$  and  $\mathcal{T}[Q_K]$
- Quiver elements can be intepreted in terms of  $T[Q_K]$  data as well as holomorphic disks:



• More details in the paper - coming soon!

(日) (周) (王) (王)

- Physically knots-quivers correspondence is a duality between 3d  $\mathcal{N} = 2$  theories  $\mathcal{T}[L_K]$  and  $\mathcal{T}[Q_K]$
- Quiver elements can be intepreted in terms of  $T[Q_K]$  data as well as holomorphic disks:



• More details in the paper - coming soon!

イロト イポト イラト イラト



# Thank you for your attention!

P. Kucharski Physics and geometry of KQ correspondence

◆□ > ◆□ > ◆豆 > ◆豆 >

nar