

# Counting instantons in $N=1$ theories of class $S_k$

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Quantum fields, knots and strings

Warsaw 25 September 2018

[1512.06079 Coman,EP,Taki,Yagi]

[1703.00736 Mitev,EP]

[1712.01288 Bourton,EP]

# Motivation: $N=2$ exact results

- \* *Seiberg-Witten theory: effective theory in the IR*
- \* *Nekrasov: instanton partition function*
- \* *Pestun: observables in the UV (path integral on the sphere localizes)*

► String/M-/F-theory realizations

- \* *Gaiotto: 4D  $N=2$  **class S**: 6D (2,0) on Riemann surface  $C_{g,n}$*
- \* *AGT: 4D partition functions = 2D CFT correlators*
- \* *4D SC Index = 2D correlation function of a TFT*

2D/4D  
relations

# What can we do for $N=1$ theories?

- ☒ *Superconformal Index* [Romelsberge 2005]  
[Kinney, Maldacena, Minwalla, Raju 2005]
- ☒ *Intriligator and Seiberg: generalized SW technology*
- ☒ *Witten: IIA/M-theory approach to curves*

Holomorphy fixes:

$N=2$  theories: prepotential (that's all in the IR)

$N=1$  theories: superpotential (there are also Kähler terms)

- ☐ *No Localization (No Nekrasov, no Pestun)*
- ☐ *An  $S^4$  partition function plagued with scheme ambiguities.*  
[Gerchkovitz, Gomis, Komargodski 2014]
- ☒ *Derivatives of the free energy scheme independent.*  
[Bobev, Elvang, Kol, Olson, Pufu 2014]

# What can we do for N=1 theories?

- ☑ Can construct **conformal** N=1 theories. [Leigh, Strassler 1995]
- ☑ AdS/CFT natural route to several examples. [Kachru, Silverstein 1998]  
[Lawrence, Nekrasov, Vafa 1998]
- ☑ 6D (1,0) on a Riemann Surface. [Gaiotto, Razamat 2015] [Heckman, Vafa....]

**Class  $S_k$  ( $S_\Gamma$ ):**

[Gaiotto, Razamat 2015]

2D/4D  
relation

- \* Conformal
- \* Obtained by orbifolding N=2 (inheritance)
- \* Labeled by punctured Riemann Surface
- \* Index = 2D correlation function of a TFT

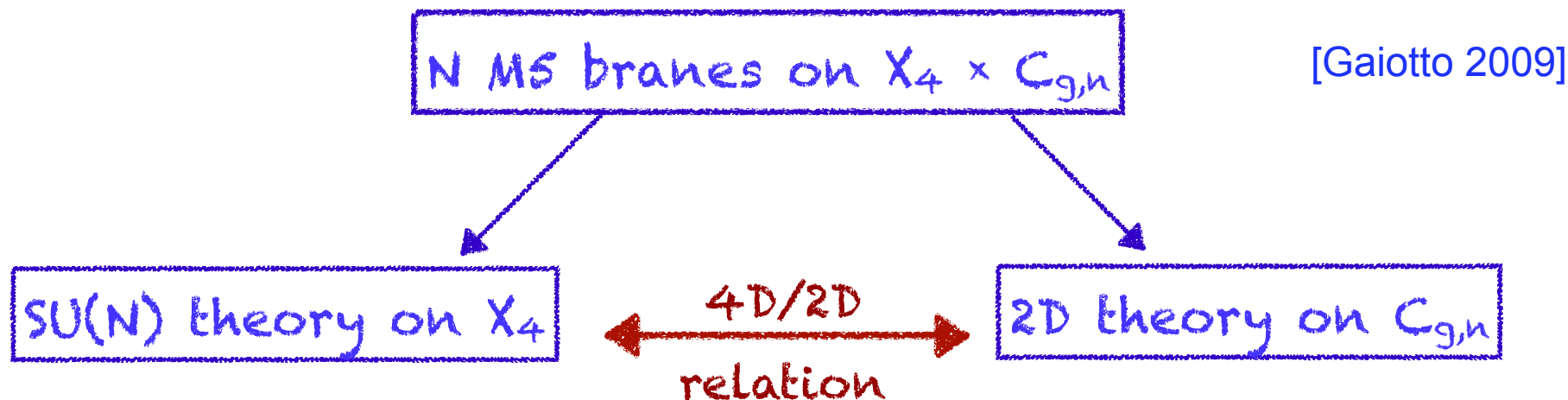
# Plan

Is there AGT<sub>k</sub>?  
4D partition functions = 2D CFT correlators

- \* *Introduce  $N=1$  theories in class  $S_k$*
- \* *Spectral curves for  $N=1$  theories in class  $S_k$*
- \* *From the curves: 2D symmetry algebra and representations*
- \* *Conformal Blocks  $\longrightarrow$  Instanton partition function*
- \* *Instanton partition function from  $Dp/D(p-4)$  branes on orbifold*
- \* *Free trinion partition functions on  $S^4 \longrightarrow$  3pt functions*

**Class  $S_k$**

# Class S and $S_k$



6D **(2,0)** SCFT on Riemann surface: 4D N=2 theories of **class S**

Transverse  $C^2/Z_k$  Orbifold the 6D (2,0) SCFT to 6D (1,0) SCFT

6D **(1,0)** SCFT on Riemann surface: 4D N=1 theories of **class  $S_k$**

	$x^0$	$x^1$	$x^2$	$x^3$	$x^4$	$x^5$	$x^6$	$x^7$	$x^8$	$x^9$	$x^{10}$
N M5-branes	—	—	—	—	.	.	—	.	.	.	—
$A_{k-1}$ orbifold	.	.	.	.	—	—	.	—	—	.	.

[Gaiotto, Razamat 2015]

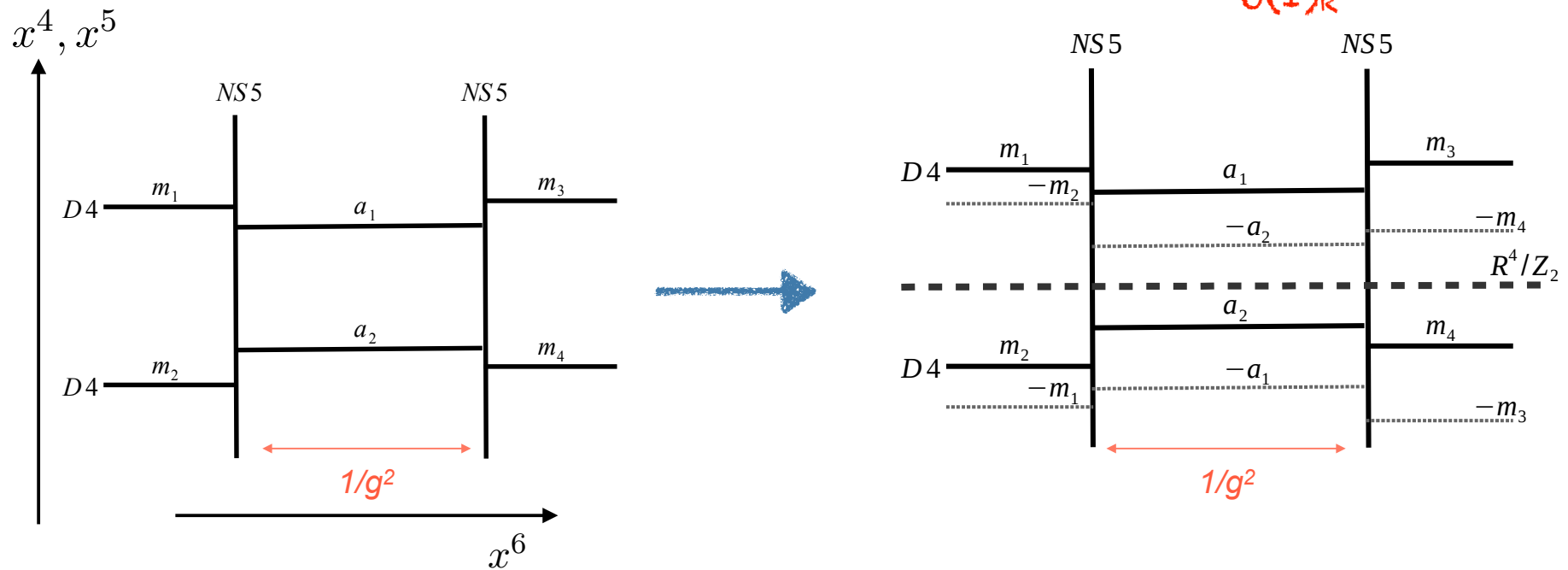
# Class $S_k$

[Gaiotto, Razamat 2015]

$U(1)_r$   ~~$SU(2)_R$~~  of  $N=2$

Type IIA

	$x^0$	$x^1$	$x^2$	$x^3$	$x^4$	$x^5$	$x^6$	$x^7$	$x^8$	$x^9$	$(x^{10})$
$M$ NS5 branes	—	—	—	—	—	—	.	.	.	.	.
$N$ D4-branes	—	—	—	—	.	.	—	.	.	.	—
$A_{k-1}$ orbifold	.	.	.	.	—	—	.	—	—	.	.



YM coupling:  $1/g^2$

$$\epsilon = \Gamma^0 \Gamma^1 \Gamma^2 \Gamma^3 \Gamma^4 \Gamma^5 \epsilon = \Gamma^0 \Gamma^1 \Gamma^2 \Gamma^3 \Gamma^6 \epsilon = \Gamma^4 \Gamma^5 \Gamma^7 \Gamma^8 \epsilon$$



# Class $S_k$

*Type IIB*

	$x^0$	$x^1$	$x^2$	$x^3$	$x^4$	$x^5$	$x^6$	$x^7$	$x^8$	$x^9$
$A_{M-1}$ orbifold	.	.	.	.	.	.	—	—	—	—
$N$ D3-branes	—	—	—	—	.	.	.	.	.	.
$A_{k-1}$ orbifold	.	.	.	.	—	—	.	—	—	.

☑  $N=1$  orbifold daughter of  $N=4$  SYM  $\Gamma = \mathbb{Z}_k \times \mathbb{Z}_M$

☑ Useful for AdS/CFT (**orbifold inheritance**)  $AdS_5 \times S^5 / (\mathbb{Z}_k \times \mathbb{Z}_M)$   
[Bershadsky, Kakushadze, Vafa 1998]

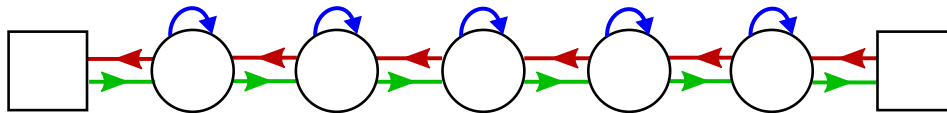
☑ String theory technics to calculate instantons

[Dorey, Hollowood, Khoze, Mattis,...] [Lerda,...]

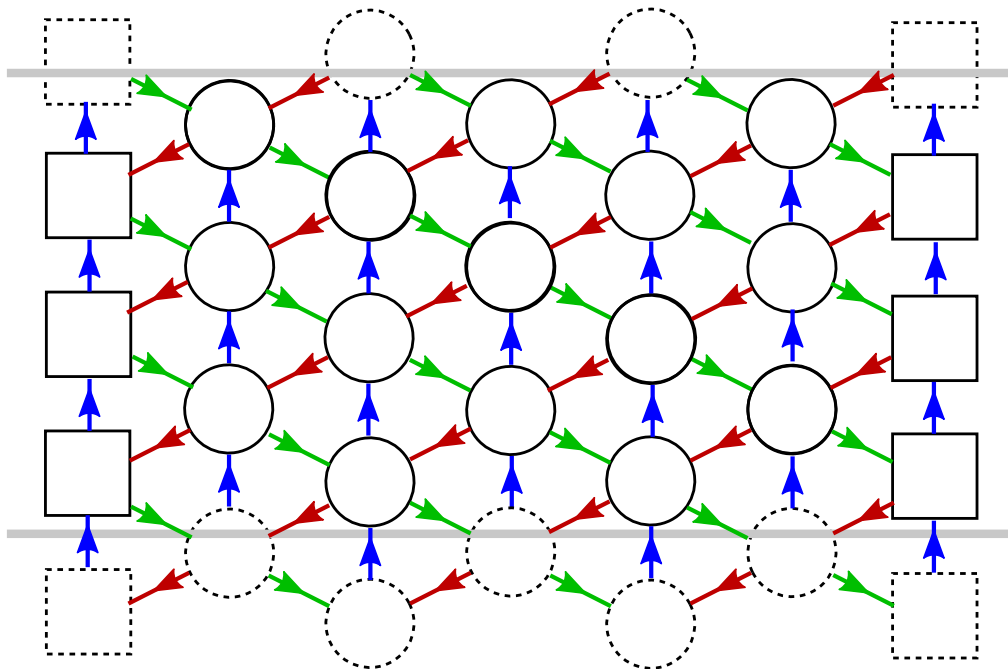
# Class $S_k$

[Gaiotto, Razamat 2015]

*4D field theory point of view*



$N=2$  class S mother theory



$N=1$  class  $S_k$  daughter theory

	$U(1)_t$	$U(1)_{\alpha_c}$	$U(1)_{\beta_{i+1-c}}$	$U(1)_{\gamma_i}$
$V_{(i,c)}$	0	0	0	0
$\Phi_{(i,c)}$	-1	0	-1	+1
$Q_{(i,c-1)}$	+1/2	-1	+1	0
$\tilde{Q}_{(i,c-1)}$	+1/2	+1	0	-1

Large global  
symmetry group

# Class $S_k$

Begin with  $N=2$  class  $S$  with  $SU(kN)$  gauge groups:

$$W_S = \sum_{c=1}^{M-1} \left( Q_{(c-1)} \Phi_{(c)} \tilde{Q}_{(c-1)} - \tilde{Q}_{(c)} \Phi_{(c)} Q_{(c)} \right)$$

$N=2$  class  $S$  mother theory

Orbifold projection: [Douglas, Moore 1996]

$$\Phi_{(c)} = \begin{pmatrix} & & & & \Phi_{(1,c)} \\ & & & & & \\ & & & & & \Phi_{(2,c)} \\ & & & & & & \ddots \\ & & & & & & & \Phi_{(k-1,c)} \\ & & & & & & & & \Phi_{(k,c)} \end{pmatrix}$$

$kN \times kN$   $N \times N$

$$Q_{(c)} = \begin{pmatrix} Q_{(1,c)} & & & \\ & Q_{(2,c)} & & \\ & & \ddots & \\ & & & Q_{(k,c)} \end{pmatrix}$$

$$\tilde{Q}_{(c)} = \begin{pmatrix} & & & \tilde{Q}_{(k,c)} \\ & & & \\ \tilde{Q}_{(1,c)} & & & \\ & \ddots & & \\ & & \tilde{Q}_{(k-1,c)} & \end{pmatrix}$$

$$W_{S_k} = \sum_{i=1}^k \sum_{c=1}^{M-1} \left( Q_{(i,c-1)} \Phi_{(i,c)} \tilde{Q}_{(i,c-1)} - \tilde{Q}_{(i,c)} \Phi_{(i,c)} Q_{(i+1,c)} \right)$$

$N=1$  class  $S_k$   
daughter theory

# Coulomb and Higgs branch

N=2 class S mother theory

\* **Coulomb Branch:**  $\langle Q \rangle = 0$  with  $\langle \phi \rangle = \text{diag}(a_1, \dots, a_N)$  and  $u_\ell = \langle \text{tr} \phi^\ell \rangle$   
 $E = r$

\* **Higgs Branch:**  $\langle \phi \rangle = a = 0$  with operators  $E = 2R$  e.g.  
 $m_i = 0$

$$\mu_{\mathcal{I}\mathcal{J}} = \langle \text{tr} (Q_{\{\mathcal{I}\}} \bar{Q}_{\mathcal{J}}) \rangle$$

N=1 class  $S_k$  daughter theory

\* **Coulomb Branch:**  $\langle Q \rangle = 0$  parameterised by  $u_{\ell k} = \langle \text{tr} (\Phi_{(1)} \cdots \Phi_{(k)})^\ell \rangle$

$$\langle U^{-1} \Phi U \rangle = \text{diag}(a_1, a_2, \dots, a_N) \otimes \text{diag} \left( 1, e^{\frac{2\pi i}{k}}, e^{\frac{4\pi i}{k}} \cdots e^{\frac{2\pi i(k-1)}{k}} \right) \quad E = r$$

\* **Higgs Branch:** similar w/ mother theory, operators charged under new beta and gamma symmetries.

**CB and HB do not mix (no relations): charged under different charges!**

# Curves

# Generic N=1 Curves

\* *The spectral curve computes the effective YM coupling constants.*

[Intriligator, Seiberg]

\* *M-theory on  $\mathbb{R}^{3,1} \times CY_3 \times \mathbb{R}^1$*  [Bah, Beem, Bobev, Wecht]

*$CY_3$  locally two holomorphic line bundles on the curve  $C_{g,n}$*

\* *N=1 spectral curve is an overdetermined algebraic system of eqns.*

[Bonelli, Giacomelli, Maruyoshi, Tanzini]

[Xie... ]

\* *For class  $S_k$  on the Coulomb Branch (  $\langle Q \rangle = 0$  ) only one equation*

*exactly like for N=2 theories.*

[Coman, EP, Taki, Yagi]

# $S_k$ curves

[Coman,EP,Taki,Yagi]

$$v = x^4 + ix^5$$

$$w = x^7 + ix^8$$

$$t = e^{-\frac{x^6 + ix^{10}}{R_{10}}}$$

$U(1)_r$   ~~$SU(2)_R$~~  of  $N=2$

	$x^0$	$x^1$	$x^2$	$x^3$	$x^4$	$x^5$	$x^6$	$x^7$	$x^8$	$x^9$	$(x^{10})$
$M$ NS5 branes	—	—	—	—	—	—	.	.	.	.	.
$N$ D4-branes	—	—	—	—	.	.	—	.	.	.	—
$A_{k-1}$ orbifold	.	.	.	.	—	—	.	—	—	.	.

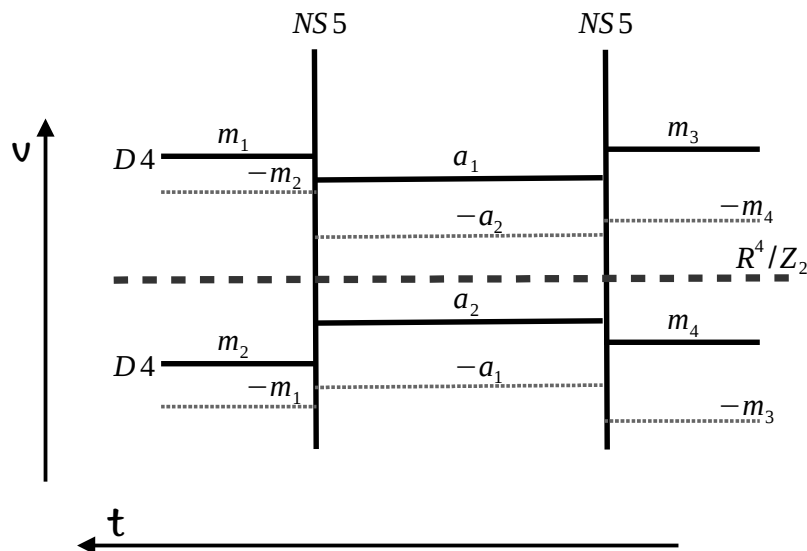
$U(1)_R$

$$(v, w) \sim \left( e^{\frac{2\pi i}{k}} v, e^{-\frac{2\pi i}{k}} w \right)$$

Take all the positions of all the branes on the  $w$  plane to be at  $w=0$   
[Lykken,Poppitz,Trivedi 97]

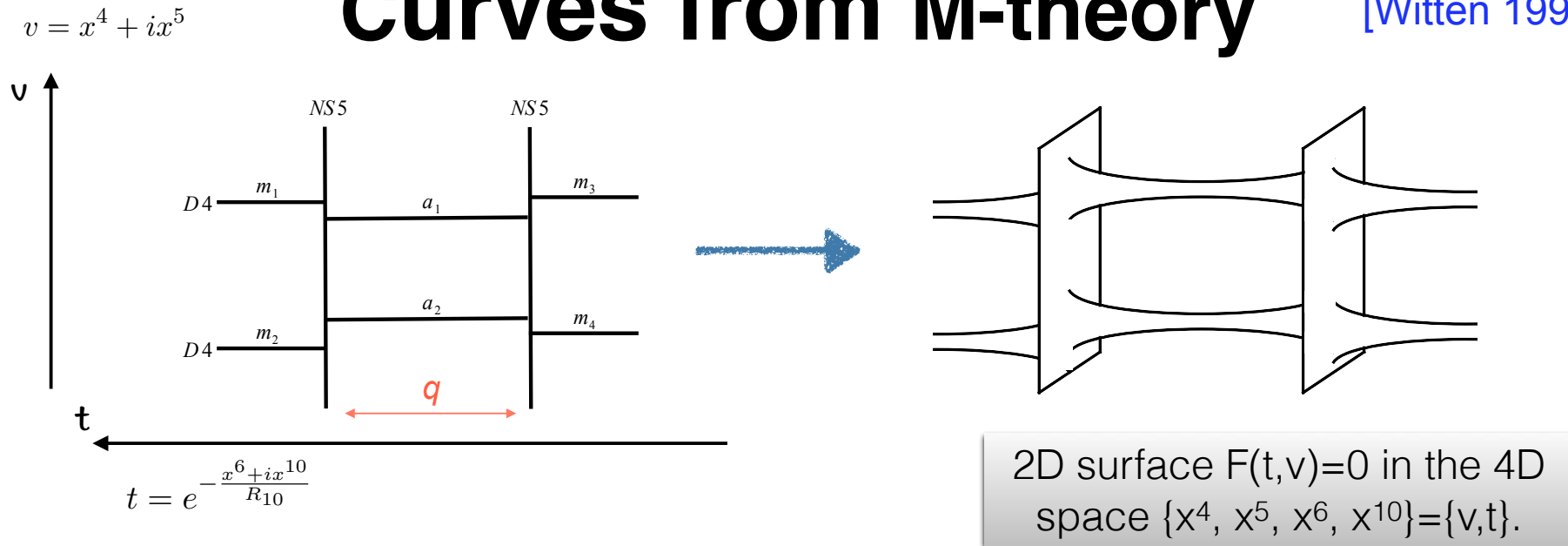
$$\langle Q \rangle = 0$$

Zero vevs for  
Higgs branch  
operators!



# Curves from M-theory

[Witten 1997]



The NS5/D4 is the classical configuration.

Take in account tension of the branes: include quantum effects.

*M-theory: a single M5 brane with non trivial topology*

$$(t - 1)(t - q)v^2 - P_1(t)v + P_2(t) = 0$$

$$(v - m_1)(v - m_2)t^2 + \left( -(1 + \underline{q})v^2 + \underline{q}\underline{M}v + \underline{u} \right)t + q(v - m_3)(v - m_4) = 0$$

coupling constant  $q = e^{2\pi i \tau w}$

$u = \text{tr} \phi^2$

$M = m_1 + m_2 + m_3 + m_4$



# Class S Curve

SU(2) with 4 flavors

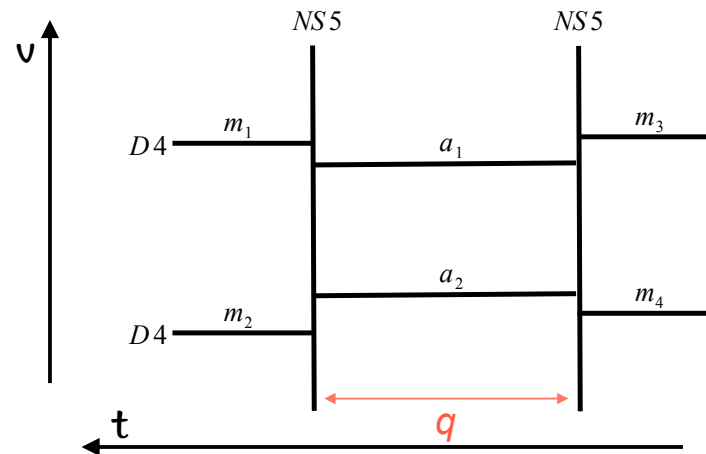
$$(t-1)(t-q)v^2 - P_1(t)v + P_2(t) = 0$$

$$(v-m_1)(v-m_2)t^2 + \left( -(1+q)v^2 + qMv + u \right) t + q(v-m_3)(v-m_4) = 0$$

$$q = e^{2\pi i \tau}$$

$$u = \text{tr} \phi^2$$

$$M = m_1 + m_2 + m_3 + m_4$$



# Class S<sub>k</sub> Curve

$$v \sim e^{\frac{2\pi i}{k}} v$$

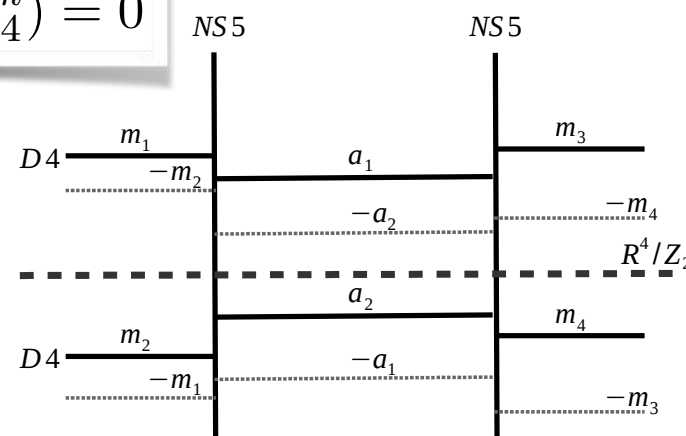
$$(v^k - m_1^k)(v^k - m_2^k)t^2 + P(v)t + q(v^k - m_3^k)(v^k - m_4^k) = 0$$

$$P(v) = -(1+q)v^{2k} + u_k v^k + u_{2k}$$

vevs of gauge invariant operators:  
parameterize the Coulomb branch

$$\langle \text{tr} (\Phi_{(1)} \cdots \Phi_{(k)}) \rangle \sim u_k$$

$$\langle \text{tr} (\Phi_{(1)} \cdots \Phi_{(k)})^2 \rangle \sim u_{2k}$$



# $S_k$ curves

SW or IR curve  $\Sigma$   
of  $g=kN-1$

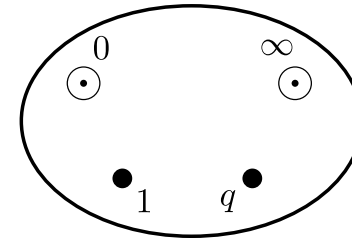


Gaiotto or UV curve  $C_{0,n}$   
a sphere with  $n$  punctures

$$x^{kN} = - \sum_{\ell=1}^N \phi_{k\ell}^{(4)}(t) x^{k(N-\ell)}$$

$$\phi_{k\ell}^{(4)}(t) = \frac{(-1)^\ell \mathfrak{c}_L^{(\ell,k)} t^2 + u_{k\ell} t + (-1)^\ell \mathfrak{c}_R^{(\ell,k)} q}{t^{k\ell} (t-1)(t-q)}$$

$$\mathfrak{c}^{(s,k)} = \sum_{i_1 < \dots < i_s = 1}^N m_{i_1}^k \cdots m_{i_s}^k$$

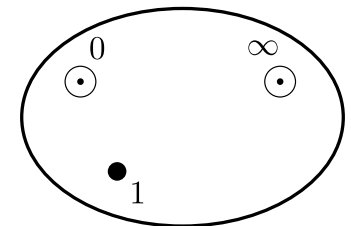


$\mathcal{N} = 1$   $SU(N)$  SCQCD $_k$

• **Full/maximal puncture:**  $k$  images of  $N$  mass parameters

• **Simple puncture:** No further parameters

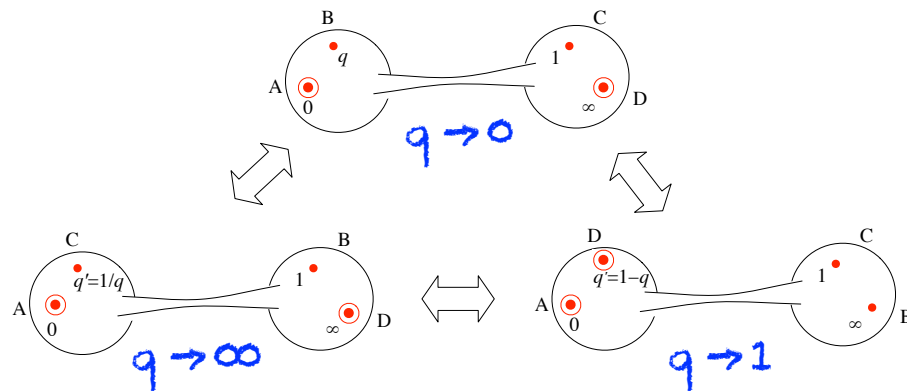
$$\phi_{k\ell}^{(3)}(t) = (-1)^\ell \frac{\mathfrak{c}_L^{(\ell,k)} t - \mathfrak{c}_R^{(\ell,k)}}{t^{k\ell} (t-1)}$$



Free Trinion  $\mathcal{C}_{0,3}^{(k)}$

# Curve decomposition and S-duality

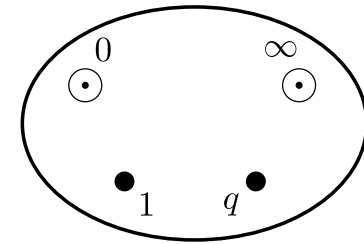
\* *Different S-duality frames (as for  $N=2$  class S)*



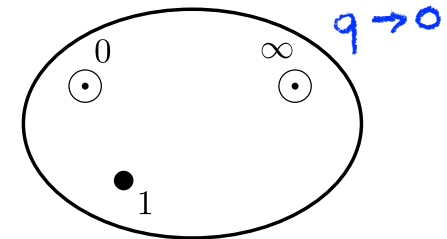
\* *Weak coupling: free trinions*

\* *Strong coupling limit: new type of punctures!*

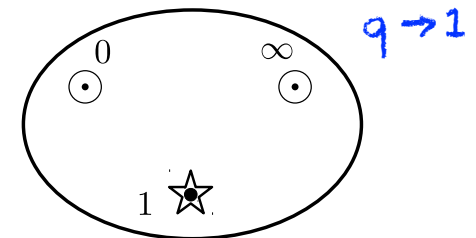
S-duality = 2D crossing equation



$\mathcal{N} = 1$   $SU(N)$  SCQCD<sub>k</sub>



Free Trinion  $\mathcal{C}_{0,3}^{(k)}$



[Coman,EP,Taki,Yagi]

# Instantons from the 2D Blocks

$$\mathcal{Z}_{\text{inst}} = \mathcal{B}_w(w_1, w_2, w_3, w_4 | q)$$

# The AGT-W correspondence

[Alday, Gaiotto, Tachikawa] [Wyllard]

*A relation between:*

► 4D N=2 theories of class S with SU(2)/SU(N) factors

► 2D Liouville/Toda CFT

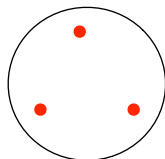
$$\mathcal{Z}_{\mathbb{S}^4} [\mathcal{T}_{g,n}] = \int da \mathcal{Z}_{pert} |\mathcal{Z}_{inst}|^2 = \int d\alpha C \dots C |\mathcal{B}_{\alpha}^{\alpha_i}|^2 = \langle \prod_{i=1}^n V_{\alpha_i} \rangle \mathcal{C}_{g,n}$$

4D gauge theory	2D CFT
instanton partition function	conformal block
perturbative part	3-point function
coupling constants	cross ratios
masses	external momenta
Coulomb moduli	internal momenta
generalized S-duality	crossing symmetry
Omega background	Coupling constant/central charge

$$b^2 = \frac{\epsilon_1}{\epsilon_2}$$

# The AGT relation from the curve

**Example:**  
 **$N=2$   $SU(2)$  Free trinion**



Close to the  
punctures:

$$\phi_2^{(3)}(z) \sim \frac{m_j^2}{(z - z_j)^2}$$

Recall 2D Ward Identities:

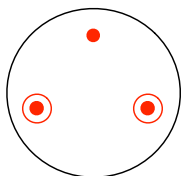
$$\langle T(z) V_1(z_1) V_2(z_2) V_3(z_3) \rangle = \sum_{j=1}^3 \left[ \frac{h_j}{(z - z_j)^2} + \frac{\partial_j}{z - z_j} \right] \langle V_1(z_1) V_2(z_2) V_3(z_3) \rangle$$

$$h_i = -m_i^2$$

$$\phi_2^{(3)}(z) = \frac{\langle T(z) V_1(z_1) V_2(z_2) V_3(z_3) \rangle}{\langle V_1(z_1) V_2(z_2) V_3(z_3) \rangle}$$

Free trinions curves are equivalent to Ward identities!

**$N=2$   $SU(N)$**   
**Free trinion**



$$\phi_\ell^{(3)}(z) = \frac{\langle W_\ell(z) V_\odot(z_1) V_\bullet(z_2) V_\odot(z_3) \rangle}{\langle V_\odot(z_1) V_\bullet(z_2) V_\odot(z_3) \rangle}$$

# From the curves to the 2D CFT

$$\lim_{\epsilon_{1,2} \rightarrow 0} \langle\langle J_\ell(t) \rangle\rangle_n = \phi_\ell^{(n)}(t)$$

$$\langle\langle J(t) \rangle\rangle_n \stackrel{\text{def}}{=} \frac{n\text{-point W-block with insertion of } J(t)}{n\text{-point W-block}}$$

- \* *The symmetry algebra that underlies the 2D CFT =  $W_{kN}$  algebra*
- \* *The reps are **standard** reps of the  $W_{kN}$  algebra*
- \* *Obtain them from the  $N=2$   $SU(kN)$  after replacing:*

$$m_{j+N_s}^{\text{SU}(Nk)} \longmapsto m_j e^{\frac{2\pi i}{k}s} \qquad a_{j+N_s}^{\text{SU}(Nk)} \longmapsto a_j e^{\frac{2\pi i}{k}s}$$

# 2D Conformal Blocks = Instanton P.F.

- \* *We have the reps of the  $W_{kN}$  algebra for  $\varepsilon_{1,2} = 0$  (from the curve)*
- \* *Demand: the structure of the multiplet (null states) not change  $\varepsilon_{1,2} \neq 0$*
- \* *The blocks for  $\varepsilon_{1,2} \neq 0$ : proposal for the instanton partition functions:*

$$\mathcal{Z}_{\text{inst}} = \mathcal{B}_{\mathbf{w}}(\mathbf{w}_1, \mathbf{w}_2, \mathbf{w}_3, \mathbf{w}_4 | q)$$

- \* *If  $w$  and  $c$  turn on  $Q \neq 0$  as in Liouville/Toda,  
then we obtain them from the  $N=2$   $SU(kN)$  after replacing:*

$$m_{j+N}^{\text{SU}(Nk)} \longmapsto m_j e^{\frac{2\pi i}{k}s} \qquad a_{j+N}^{\text{SU}(Nk)} \longmapsto a_j e^{\frac{2\pi i}{k}s}$$



# Instantons from $D(p-4)$ branes

# Instantons from D branes

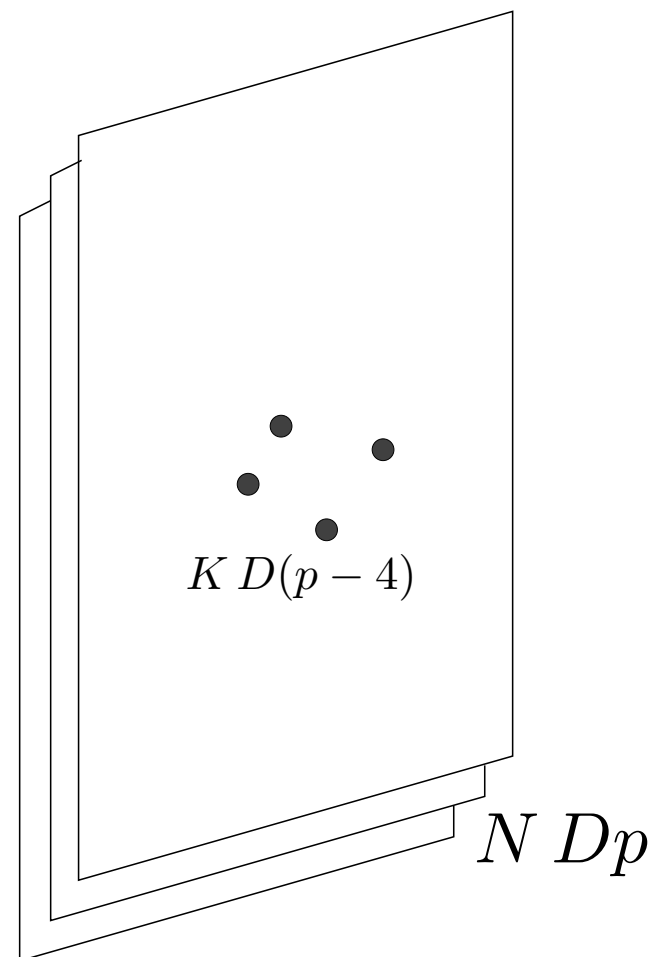
$$\text{Tr} \int_{Dp} d^{p+1}x \quad C_{p-3} \wedge F \wedge F$$

**Instanton on Dp brane = D(p-4) brane**

$$\mu^{\mathbf{C}} := [B_1, B_2] + IJ = 0 \ ,$$

$$\mu^{\mathbf{R}} := [B_1, B_1^\dagger] + [B_2, B_2^\dagger] + II^\dagger - J^\dagger J = 0$$

The Dp-D(p-4) strings give the N x K I, J<sup>†</sup>  
and the D(p-4)-D(p-4) the K x K B<sub>1</sub>, B<sub>2</sub>  
auxiliary matrices of the ADHM  
construction



$$\mathcal{M}_{K\text{-inst}}^{Dp} \simeq \mathcal{M}_{\text{Higgs}}^{D(p-4)}$$

$$= \{B_1, B_2, I, J \mid \mu^{\mathbf{C}} = 0, \mu^{\mathbf{R}} = 0\} / U(K)$$

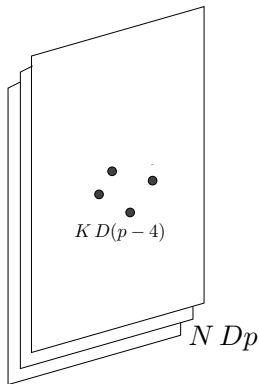
$$= \{X_{Dp}^\perp = 0, \mathcal{V}_{D(p-4)} = 0\} / U(K)$$

[Witten 1995, Douglas 1995, Dorey 1999, ...]

# Instanton Partition Function from 2D SCI

$$\mathcal{Z}_{K\text{-inst}}(a, m, \epsilon_1, \epsilon_2) = \text{Tr}_{\mathcal{M}_{K\text{-inst}}} \underline{(-1)^F e^{\epsilon_1 J_L} e^{\epsilon_2 J_R} e^{\xi R} e^{a J_G} e^{m J_F}}$$

Nekrasov's Instantons: deformed Witten index



For  $p=5$  and when the D1s wrapping a  $T^2$ :

$$\mathcal{M}_{K\text{-inst}}^{Dp} \simeq \mathcal{M}_{\text{Higgs}}^{D(p-4)}$$

Instanton's index =  
2D SCI = flavoured elliptic genus

of the 2D gauge theory living on the D1s

$$\mathcal{Z}_{K\text{-inst}}^{D5}(a, m, \epsilon_1, \epsilon_2) = \mathcal{Z}_{\text{Higgs}}^{K D1}(a, m, \epsilon_1, \epsilon_2) = \mathcal{I}_{2D} = \text{Tr}_{\mathcal{M}_{\text{Higgs}}^{K D1}} (-1)^F e^{\epsilon_1 J_L} e^{\epsilon_2 J_R} e^{\xi R} e^{a J_G} e^{m J_F}$$

2D R-symmetry  $J_{01}, J_{23}$   
 $J_{69} - J_{78}$   
 $J_{69} + J_{78}$  2D Flavor

2D SCI is very well studied, very easy to compute!

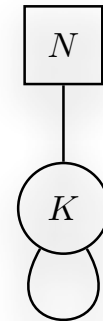
[Gadde, Gukov, Putrov]  
[Benini, Eager, Hori, Tachikawa]

# Mass deformed N=4 SYM

First we practice without the orbifolds:

- \* Obtain the  $\mathcal{Z}_{\text{inst}}$  for mass deformed N=4 SYM (N=2\*)  
from a **2D SCl computation**: 2D theory with (4,4) susy

$$D3/D(-1) \xleftrightarrow{\text{T-duality}} D5/D1$$



- \* The 2D SCl of the gauge theory on the D1 branes  
= instantons of the 6D (2,0) on  $\mathbb{R}^4 \times T^2$  = M-strings.
- \* KK reducing: instantons of mass deformed 5D N=2 MSYM on  $S^1$   
and further KK reduce to mass deformed N=4 SYM in 4D.

# Instantons with an orbifold

- \* **“Orbifold”** the 2D SCI of mass deformed  $N=4$  SYM with one, say the  $Z_M$  orbifold: we get  $M$ -strings on a transverse orbifold  
= instantons of an  $SU(N)^M$  quiver when reduce down to 5D/4D

2D theory with (0,4) susy

- \* Further **“Orbifold”** the 2D SCI with the  $Z_k$  orbifold we get something new that should correspond to instantons of class  $Sk$ ,  
the rational, trigonometric and elliptic uplift.

2D theory with (0,2) susy

# S<sub>k</sub> Instantons

$$\mathcal{Z}_{K\text{-inst}}^{\text{D5}}(a, m, \epsilon_1, \epsilon_2) = \mathcal{Z}_{\text{Higgs}}^{K\text{ D1}}(a, m, \epsilon_1, \epsilon_2) = \mathcal{I}_{2D} = \text{Tr}_{\mathcal{M}_{\text{Higgs}}^{K\text{ D1}}}(-1)^F e^{\epsilon_1 J_L} e^{\epsilon_2 J_R} e^{\xi R} e^{a J_G} e^{m J_F}$$

	$x^0$	$x^1$	$x^2$	$x^3$	$x^4$	$x^5$	$x^6$	$x^7$	$x^8$	$x^9$
$N\text{ D5}$	—	—	—	—	—	—	·	·	·	·
$\mathbb{Z}_M$	·	·	·	·	·	·	×	×	×	×
$\mathbb{Z}_k$	·	·	·	·	×	×	·	×	×	·
$K\text{ D1}$	·	·	·	·	—	—	·	·	·	·

$$\mathbb{R}^4 \times T^2$$

$$\mathbb{R}^4 \times S^1$$

$$\mathbb{R}^4$$

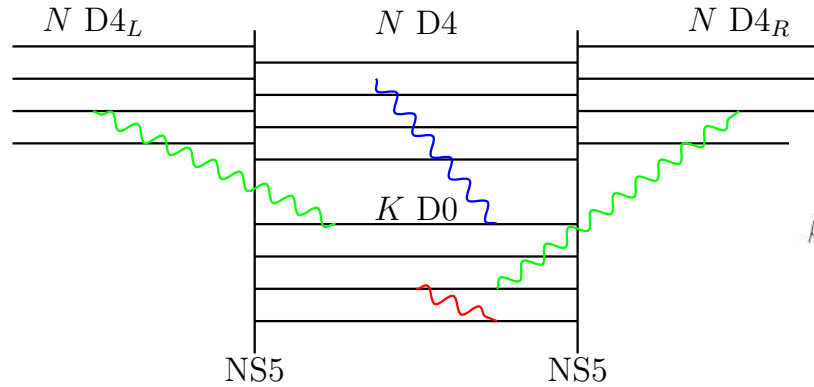
$$\mathcal{I}_{D1} = \mathcal{Z}_{\text{inst}}^{(2,0)} \xrightarrow{\text{KK}} \mathcal{Z}_{\text{inst}}^{\mathcal{N}=2} \xrightarrow{\text{KK}} \mathcal{Z}_{\text{inst}}^{\mathcal{N}=4}$$

$$\mathcal{I}_{D1}^{\mathbb{Z}_M} = \mathcal{Z}_{\text{inst}}^{(1,0)} \xrightarrow{\text{KK}} \mathcal{Z}_{\text{inst}}^{\mathcal{N}=1} \xrightarrow{\text{KK}} \mathcal{Z}_{\text{inst}}^{\mathcal{N}=2}$$

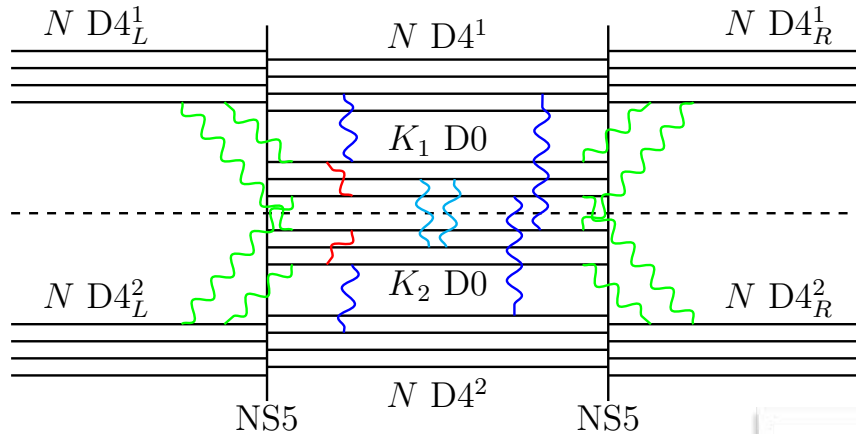
$$\mathcal{I}_{D1}^{\mathbb{Z}_M \times \mathbb{Z}_k} = \mathcal{Z}_{\text{inst}}^{(1,0)\text{w/defect}} \xrightarrow{\text{KK}} \mathcal{Z}_{\text{inst}}^{\mathcal{N}=1\text{w/defect}} \xrightarrow{\text{KK}} \boxed{\mathcal{Z}_{\text{inst}}^{\mathcal{N}=1}}$$

✱ The same answer as in [\[1703.00736 Mitev,EP\]](#) from conformal blocks!

# Instantons from D0 branes



$$\lim_{\beta_1, \beta_2 \rightarrow 0} I_{\text{SCI}} = \frac{1}{K!} \int \prod_{l=1}^K du_l \prod_{l,j=1}^K \frac{u'_{lj} (u_{lj} - 2\epsilon_+)}{(u_{lj} + \epsilon_1)(u_{lj} + \epsilon_2)} \times \prod_{l=1}^K \prod_{A=1}^N \frac{(u_l - m_{L,A})(u_l - m_{R,A})}{(u_l - a_A - \epsilon_+)(u_l - a_A + \epsilon_+)}$$



$$\lim_{\beta_1, \beta_2 \rightarrow 0} I_{\text{SCI}}^{\text{orb}} \propto \prod_{i=1}^k \int \prod_{l=1}^{K_i} du_{i,l} \prod_{l=1}^{K_i} \frac{u'_{ii,l} \prod_{j \neq i} \prod_{J=1}^{K_j} (u_{ij,lJ} - 2\epsilon_+)}{\prod_{j=1}^k \prod_{J=1}^{K_j} (u_{ij,lJ} + \epsilon_1)(u_{ij,lJ} + \epsilon_2)} \times \prod_{j=1}^k \prod_{l=1}^{K_i} \prod_{A=1}^N \frac{(u_{i,l} - \tilde{m}_{L,j,A})(u_{i,l} - \tilde{m}_{R,j,A})}{(u_{i,l} - \tilde{a}_{j,A} - \epsilon_+)(u_{i,l} - \tilde{a}_{j,A} + \epsilon_+)}$$

# **The Perturbative piece**



# Free trinion P.F. = 2D CFT 3pt functions

[in progress Carstensen,EP,Mitev]

$$Z_{\text{free trinion}}^{S^4} = \langle V_{\odot}(\infty) V_{\bullet}(1) V_{\odot}(0) \rangle$$

- \* *The free trinion theory on  $S^4$  : we can do the determinant!*
- \* *We know the conformal blocks: can write crossing equations*
- \* *Is the free trinion P.F. a solution of the crossing equations ??*
- \* *3pt functions (dynamics) + Blocks = AGT<sub>k</sub>*

# The Perturbative part

❑ *No Localization: No Pestun for  $N=1$  susy!*

\* *The free trinion theory on  $S^4$  : we can do the determinant!*

[in progress Carstensen,EP,Mitev]

\* *Guess via analytic continuation from 3D* [Gorantis,Minahan,Naseer]

*The orbifolded hyper part agrees with our determinant calculation!*

\* *Deconstruction of  $(1,1)$  Little Strings* [Hayling,Panerai,Papageorgakis]

*vector and chiral contributions precisely give the perturbative Little string!*

We can write a reasonable guess for the perturbative part!

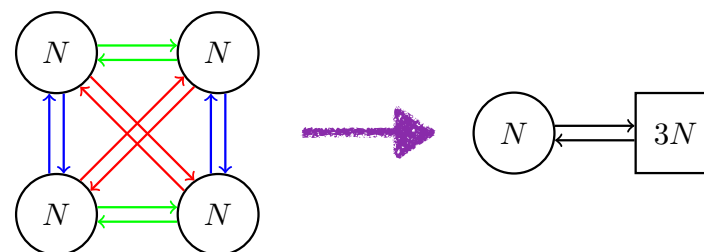
# Summary

- ☑ *We constructed spectral **curves** for  $N=1$  theories in class  $S_k$*
- ☑ *The curves: 2D symmetry algebra ( $W_{kN}$ ) and representations*
- ☑ *Conformal Blocks  $\longrightarrow$  Instanton partition function*
- ☑ *Instanton partition function from  $Dp/D(p-4)$  on orbifold*
- ☐ *Free trinion partition functions on  $S^4 =$  **3pt functions***  
[in progress Carstensen,EP,Mitev]

# Future

\* *Compute one, two instantons with standard QFT techniques*

\* *Go away from the orbifold point*  
[to appear Bourton, EP]



\* *Other  $N=1$  theories*

\* *The perturbative part?  $N=1$  partition function on  $S^4$ !*

\* *Get the  $AGT_k$  from  $(1,0)$  6D à la Cordova and Jafferis*

**Thank you!**